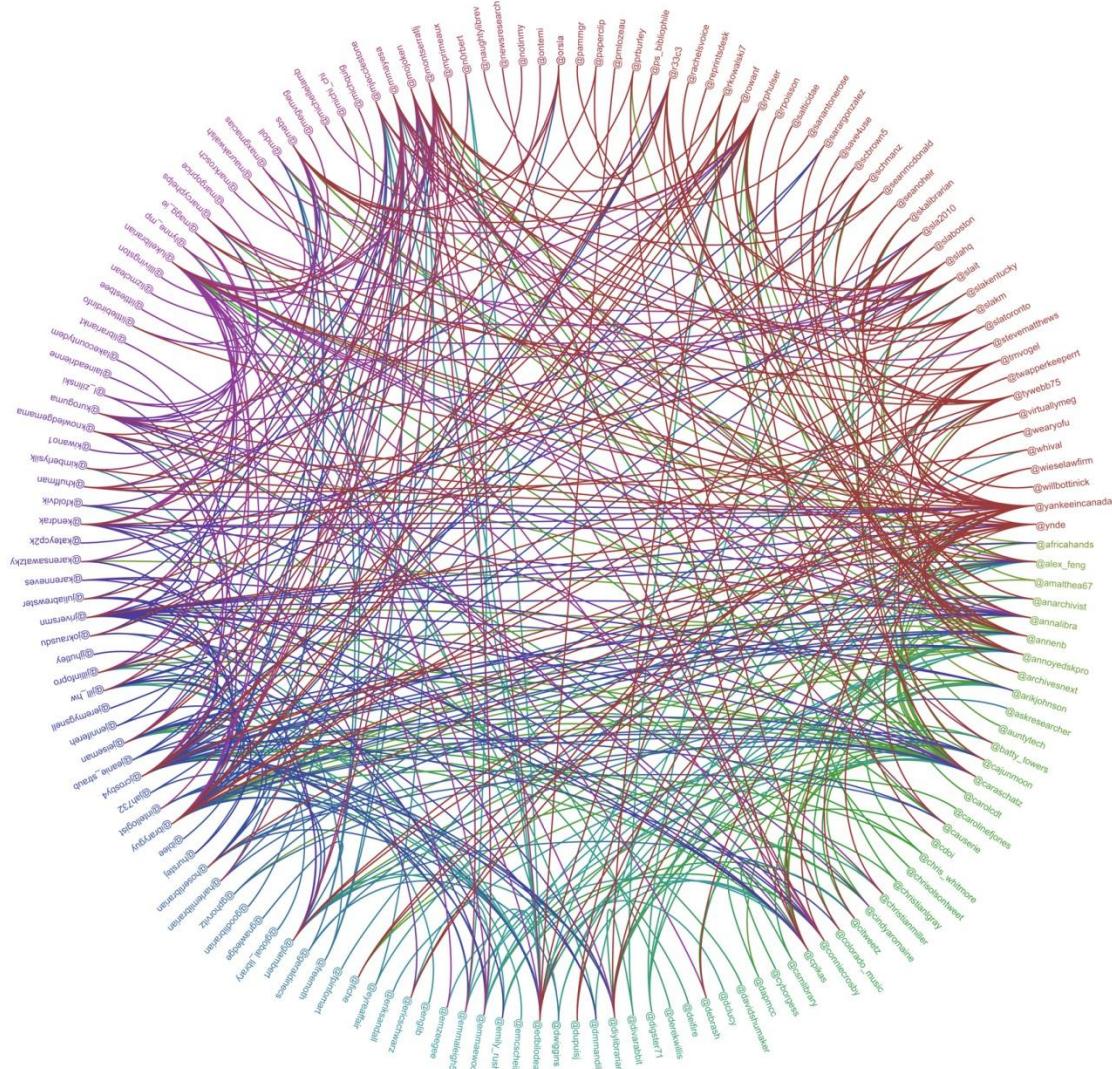


PDE Approaches to Graph Analysis



Justin Solomon
Geometric Computing Group
Stanford University

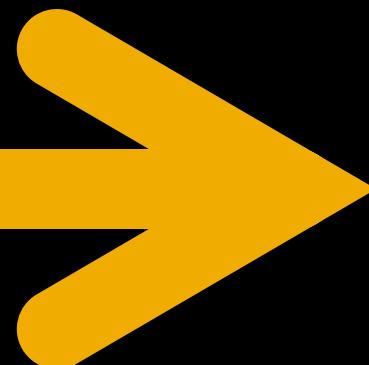
Understanding Graph Structure



Understanding Graph Structure

Attractive but
not informative.

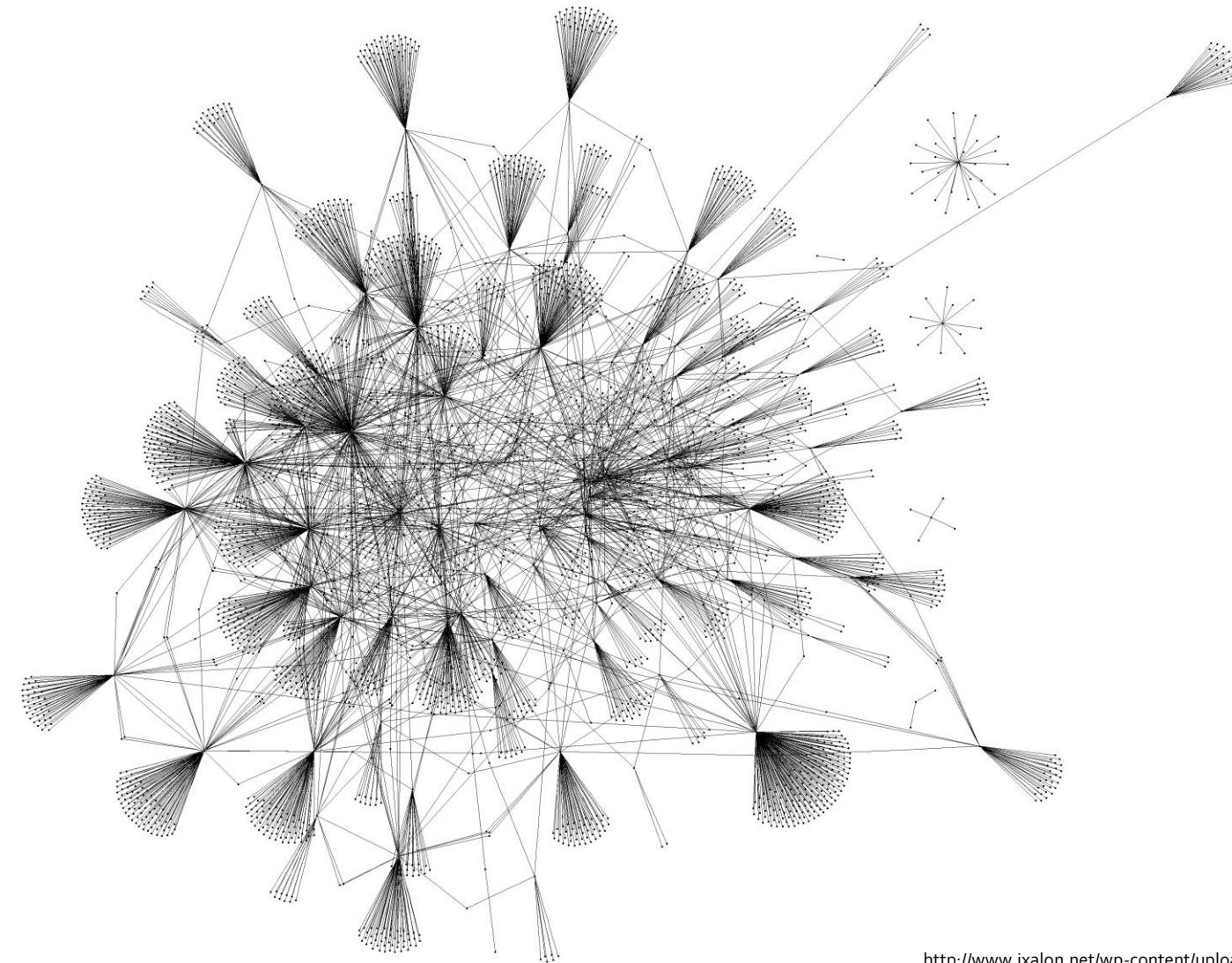
Topology [*tuh-pol-uh-jee*]:
The study and
characterization of a domain's
connectivity.



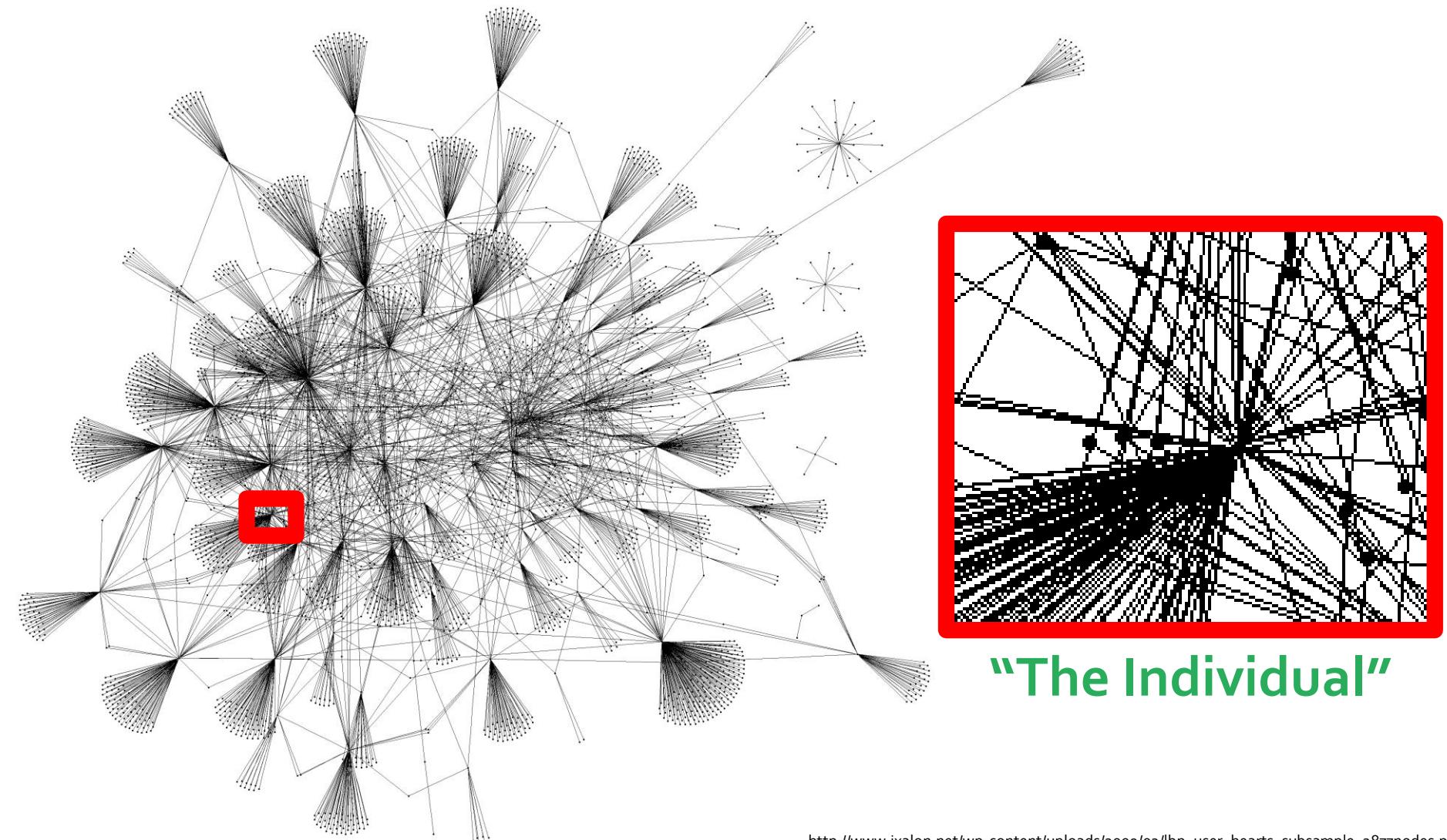
Multiscale Analysis

**Node's role changes
depending on
neighborhood**

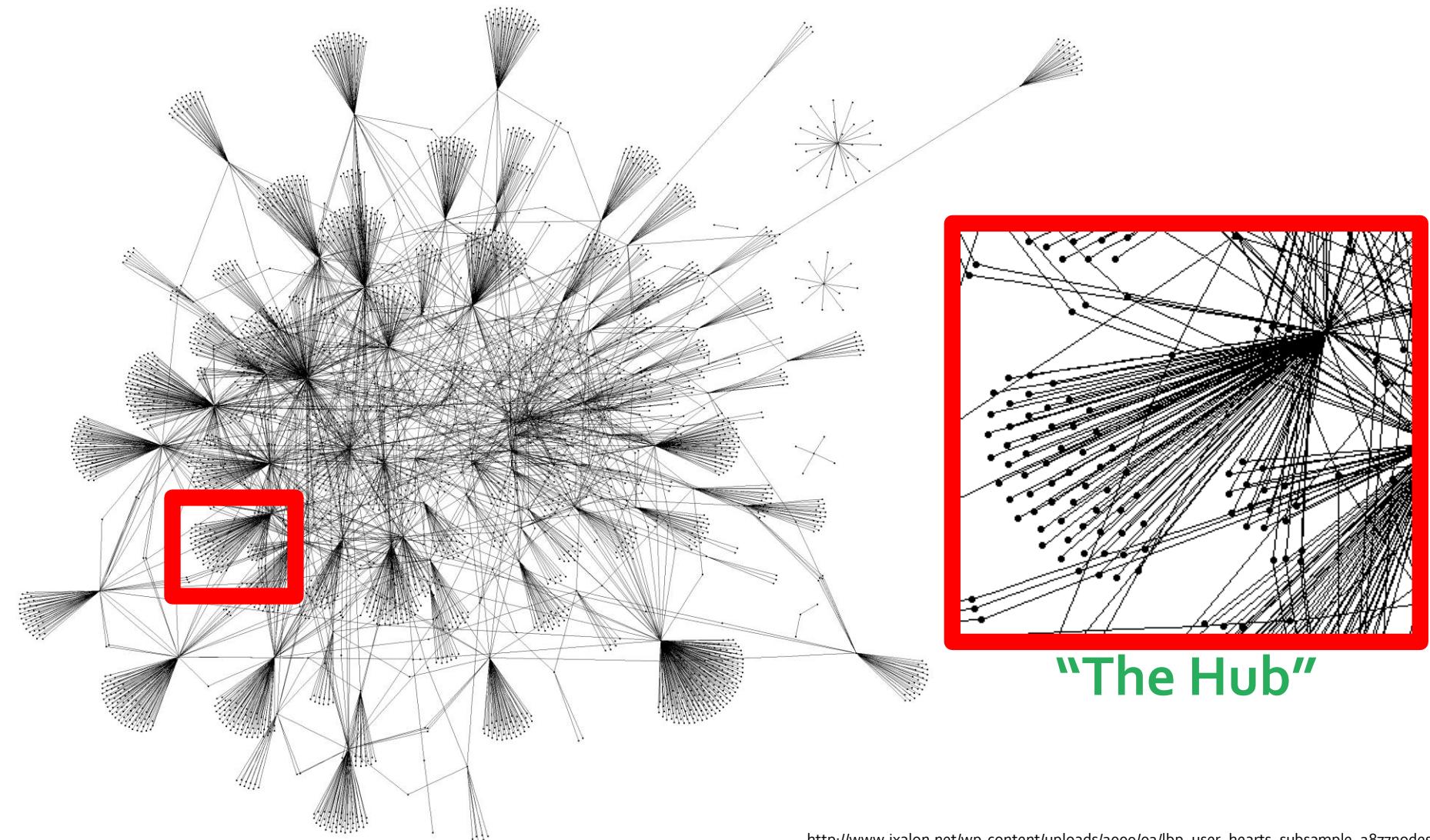
Multiscale Analysis



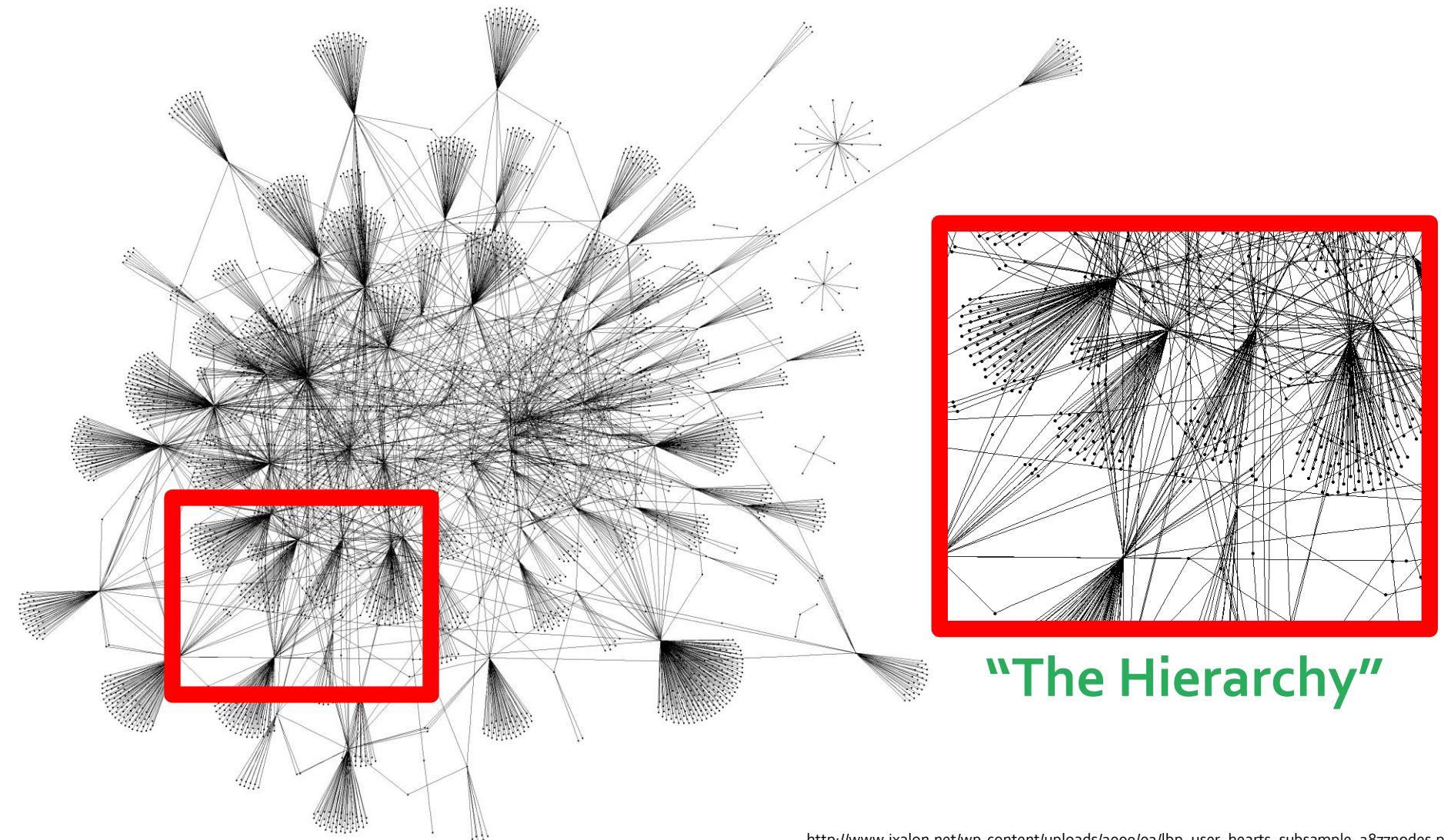
Multiscale Analysis



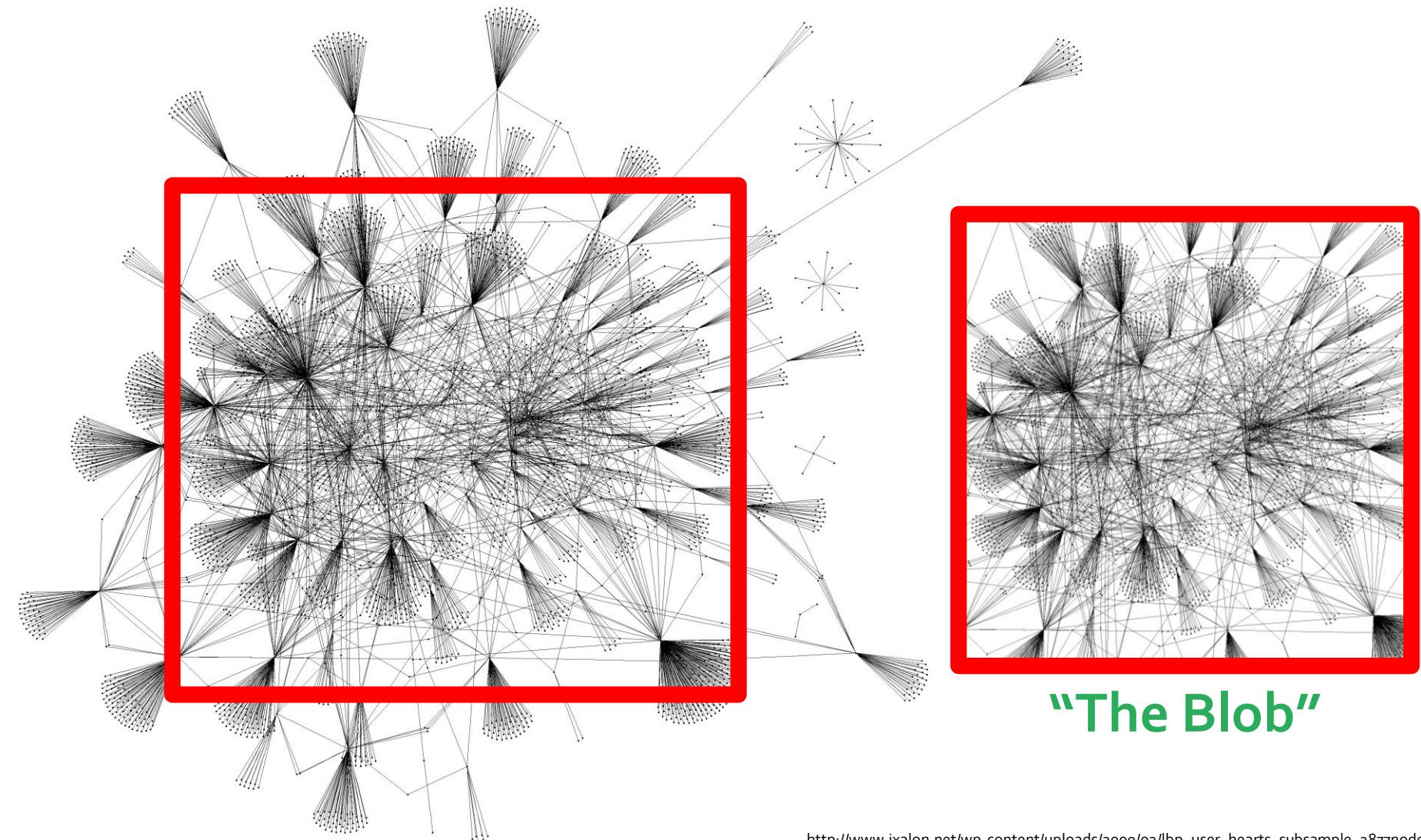
Multiscale Analysis



Multiscale Analysis



Multiscale Analysis



Potentially Interesting Function

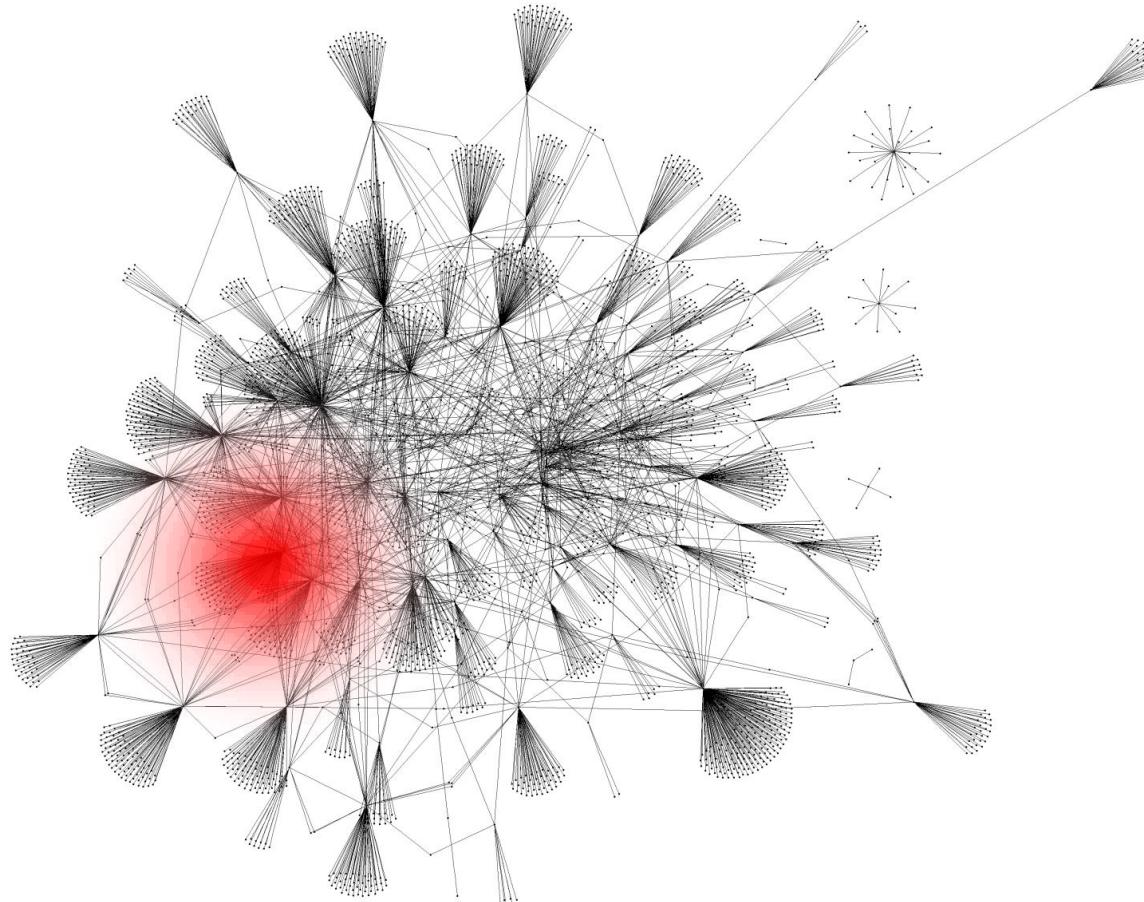
Topology(t)

Potentially Interesting Function

Topology(t)

Composed of
discrete events.

Alternative Viewpoint



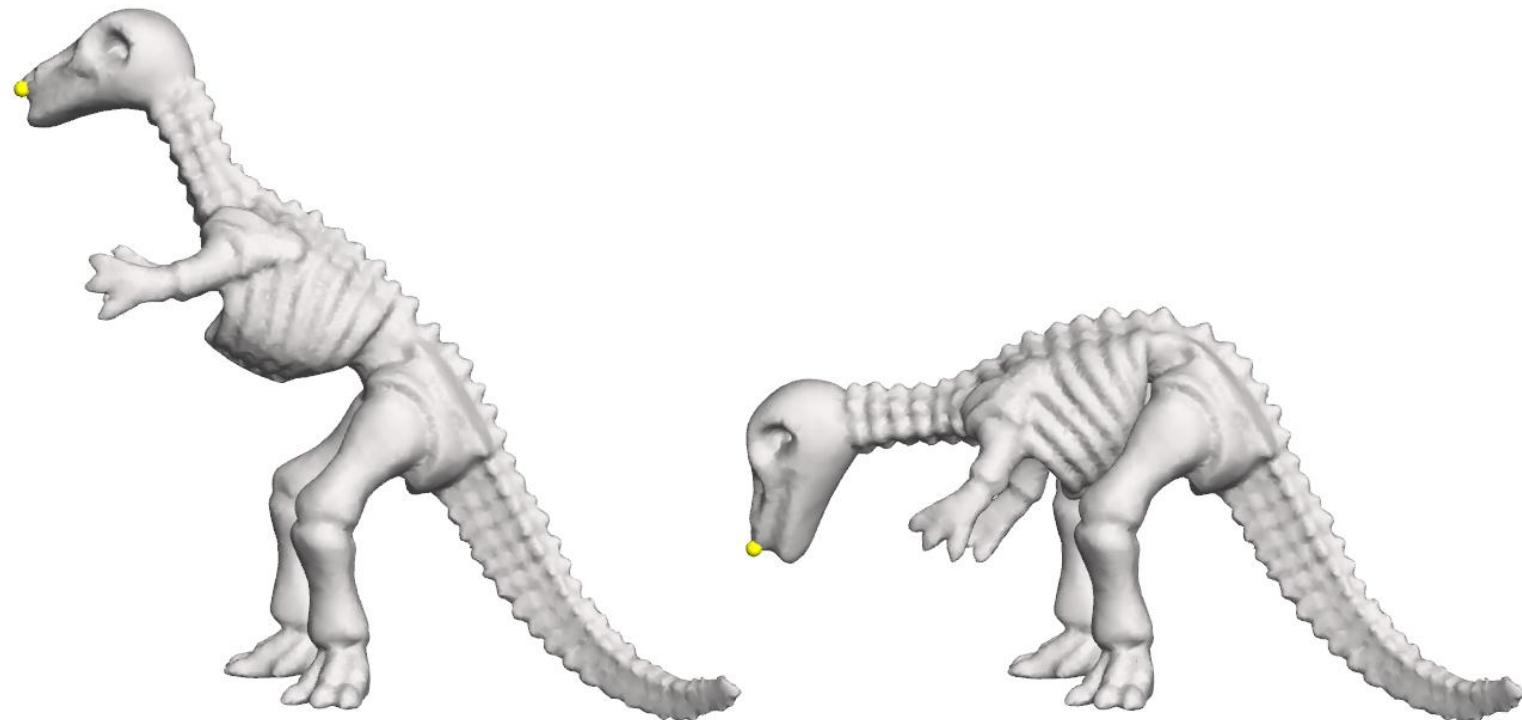
Gaussian weighting

Goal

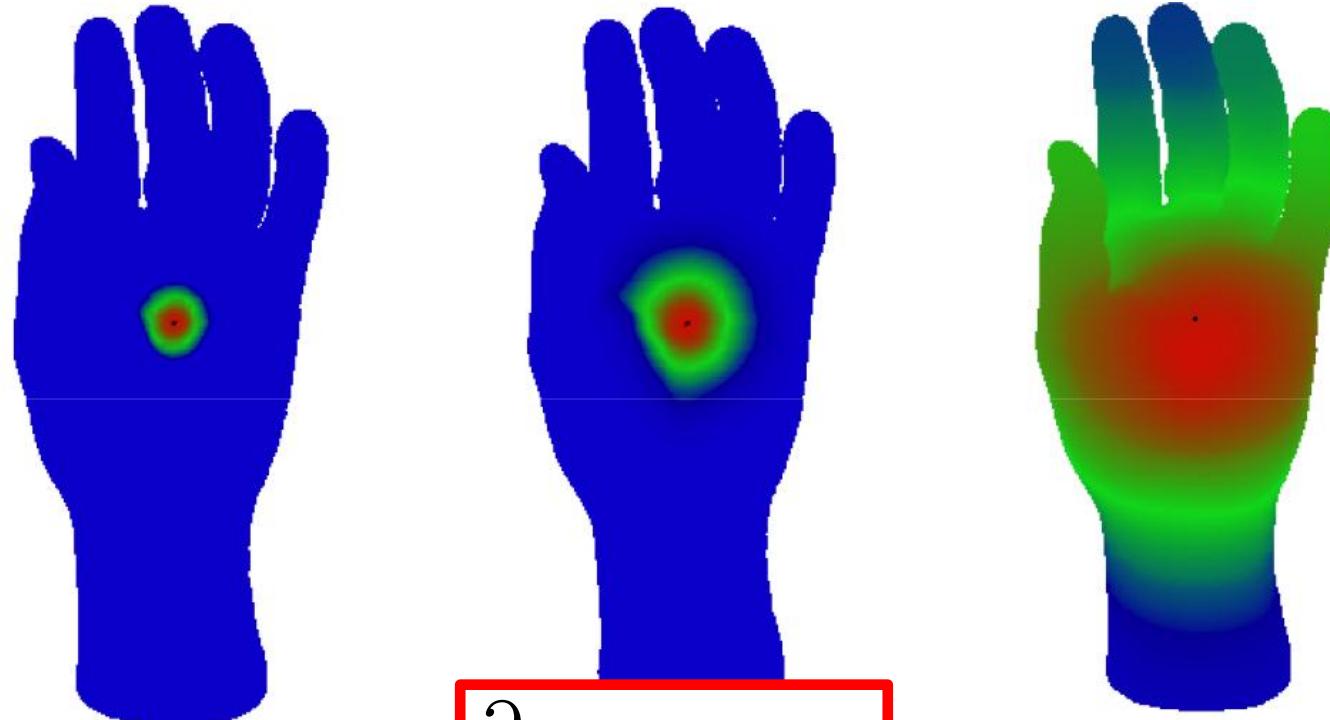
Topological analysis with
continuous dependence
on scale.

Success Story:

Descriptor-Based Matching



Success Story: Descriptor-Based Matching

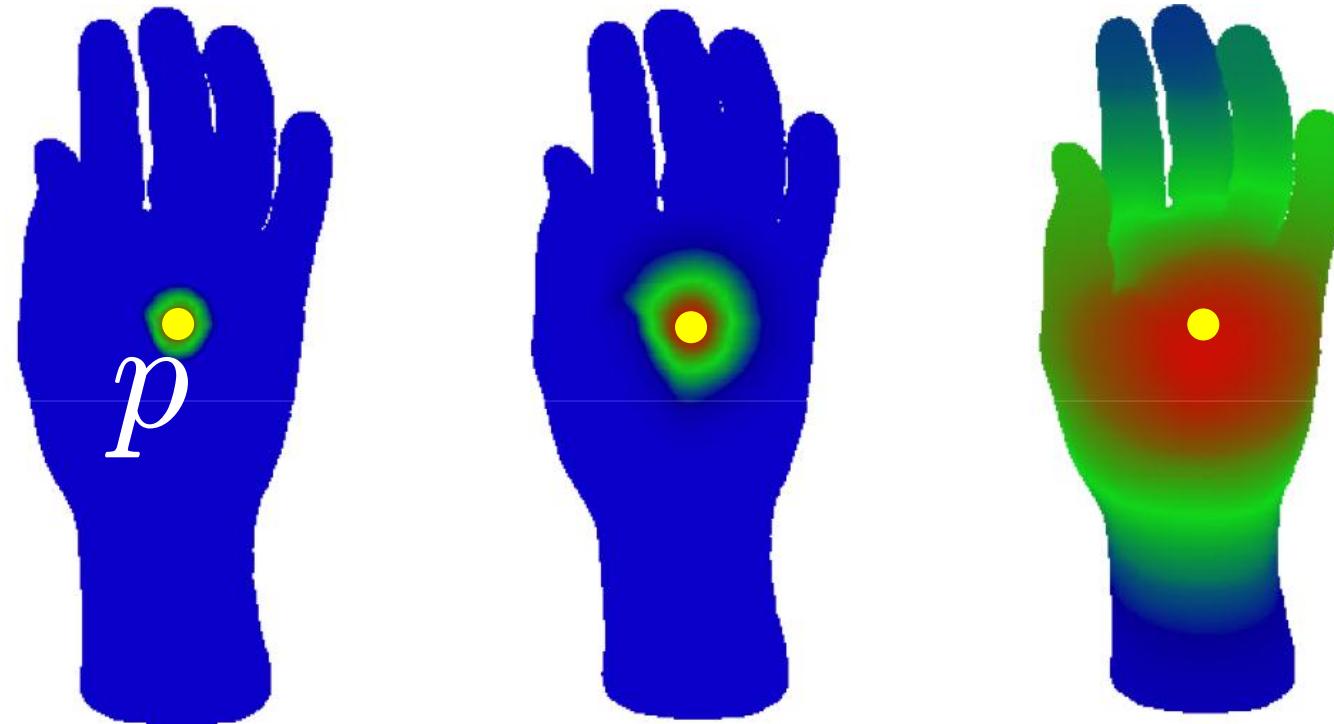


$$\frac{\partial u}{\partial t} = -\Delta u$$

http://graphics.stanford.edu/courses/cs468-10-fall/LectureSlides/11_shape_matching.pdf

Heat equation

Success Story: Descriptor-Based Matching



$HKS_p(t)$ = Heat left at time t

http://graphics.stanford.edu/courses/cs468-10-fall/LectureSlides/11_shape_matching.pdf

Heat Kernel Signature

Problem

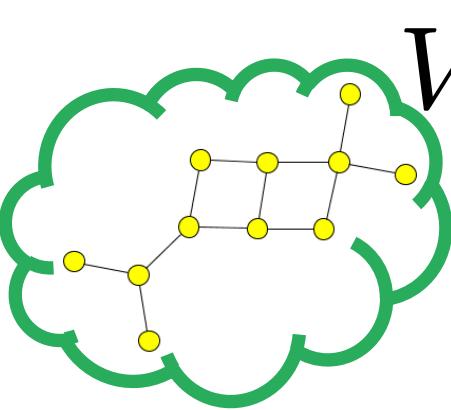
$$G = (V, E)$$

$$\begin{aligned} V &= \{v_1, v_2, v_3, \dots, v_n\} \\ E &\subseteq V \times V \end{aligned}$$

Graphs are topological objects

Problem

$$G = (V, E)$$


$$V = \{v_1, v_2, v_3, \dots, v_n\}$$
$$E \subseteq V \times V$$

<http://graphml.graphdrawing.org/primer/simple.png>

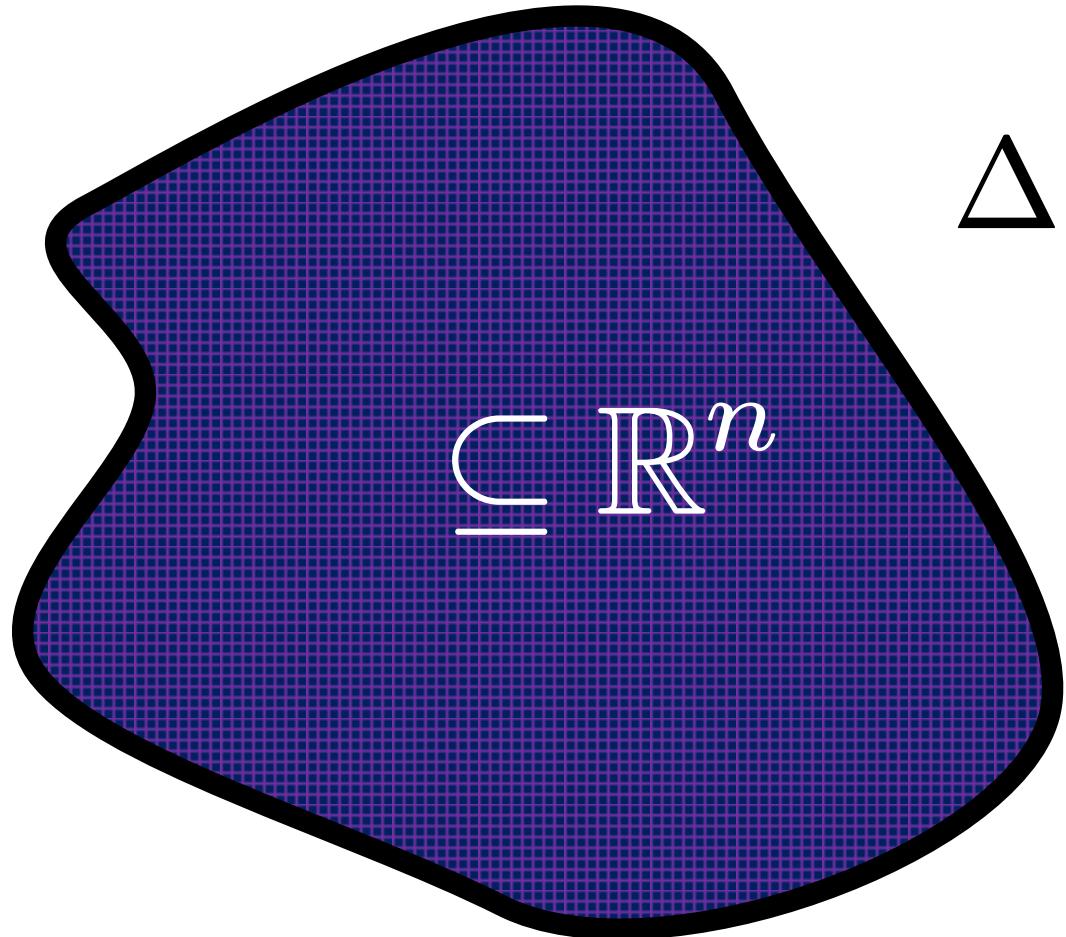
Graphs are topological objects

<question>

How do you model
flows and waves
on graphs?

</question>

The Laplacian



$$\Delta = - \sum_i \frac{\partial^2}{\partial x_i^2}$$

Key Properties

- **Linear**

$$\Delta(c_1 u + c_2 v) = c_1 \Delta u + c_2 \Delta v$$

- **Compact and bounded**

$$\|\Delta u\| \leq M \|u\| \quad \forall u$$

- **Self-adjoint**

$$\langle \Delta u, v \rangle = \langle u, \Delta v \rangle$$

- **Categorizes extrema**

$$x \in \text{local minimum} \implies [\Delta u](x) \geq 0$$

Key Properties

- **Linear**

$$\Delta(c_1 u + c_2 v) = c_1 \Delta u + c_2 \Delta v$$

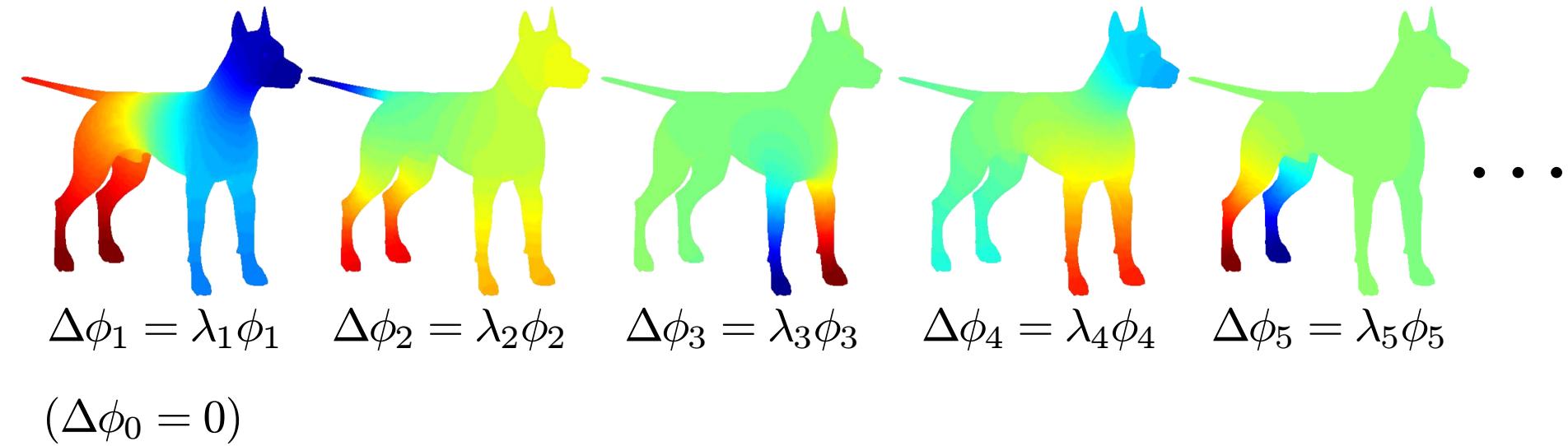
- **Compact and bounded**
- **Sufficient for many model equations.**

$$\langle \Delta u, v \rangle = \langle u, \Delta v \rangle$$

- **Categorizes extrema**

$$x \in \text{local minimum} \implies [\Delta u](x) \cdot 0$$

Laplacian Eigenfunctions



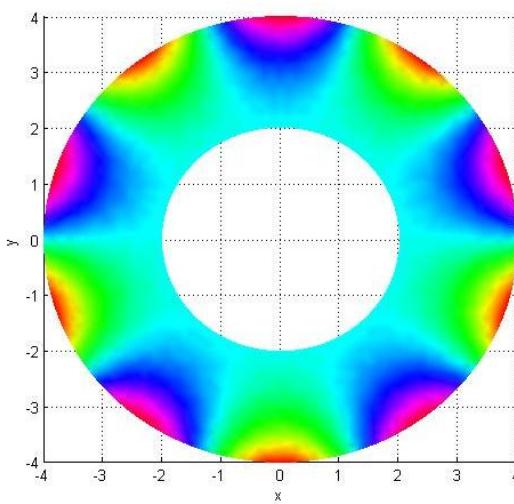
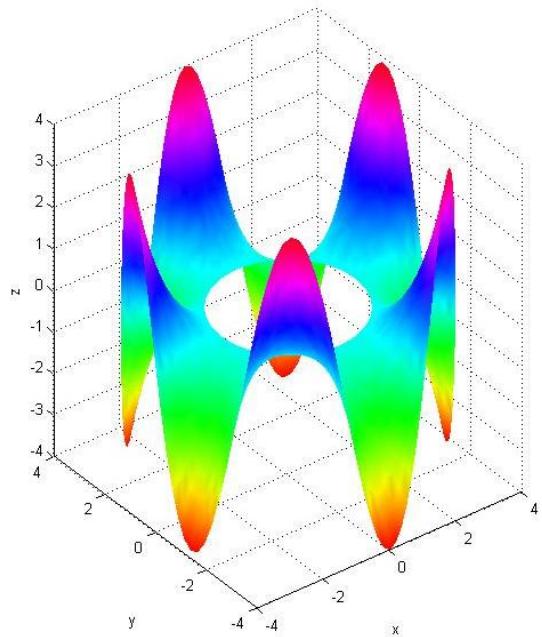
Prescribe boundary conditions on $S \subseteq \Omega$.

Analogous to Fourier basis

Model PDEs

$$\Delta u = 0$$

$$\Delta u = f$$



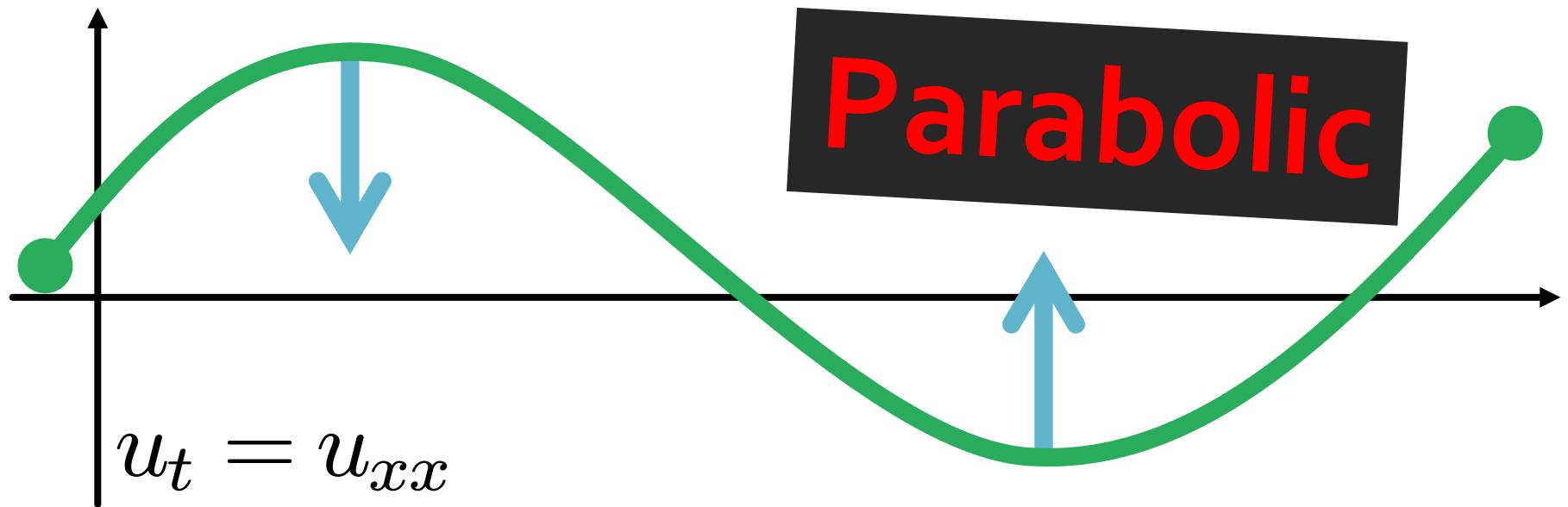
Elliptic

http://upload.wikimedia.org/wikipedia/commons/8/85/Laplace%27s_equation_on_an_annulus.jpg

Poisson and Laplace equations

Model PDEs

$$u_t = -\Delta u$$



Parabolic

Heat equation

Model PDEs

$$u_{tt} = -\Delta u$$

Hyperbolic

http://www.youtube.com/watch?v=l_yxwgh7Nbc&feature=related

Wave equation

Eigenfunction Solutions

Heat equation

$$u = \sum_i a_i e^{-\lambda_i t} \phi_i$$

Wave equation

$$u = \sum_{\lambda_i=0} (a_i + b_i t) \phi_i + \sum_{\lambda_i \neq 0} a_i \cos(\sqrt{\lambda_i} t + b_i) \phi_i$$

Generic Strategy

1. Define the domain Ω .
2. Define the Laplacian Δ .
3. See what happens.

Three Approaches

WAVE EQUATIONS FOR GRAPHS AND THE
EDGE-BASED LAPLACIAN

JOEL FRIEDMAN AND JEAN-PIERRE TILlich

In the last few years, there has made significant progress in the theory of wave equations on graphs. This note reviews some of the main results and applications, and discusses some open problems.

Discrete Green's Functions

Fan Chung¹

University of California, San Diego, La Jolla, California 92093-0112

and

S.-T. Yau

Diffusion and Elastic Equations on Networks

By

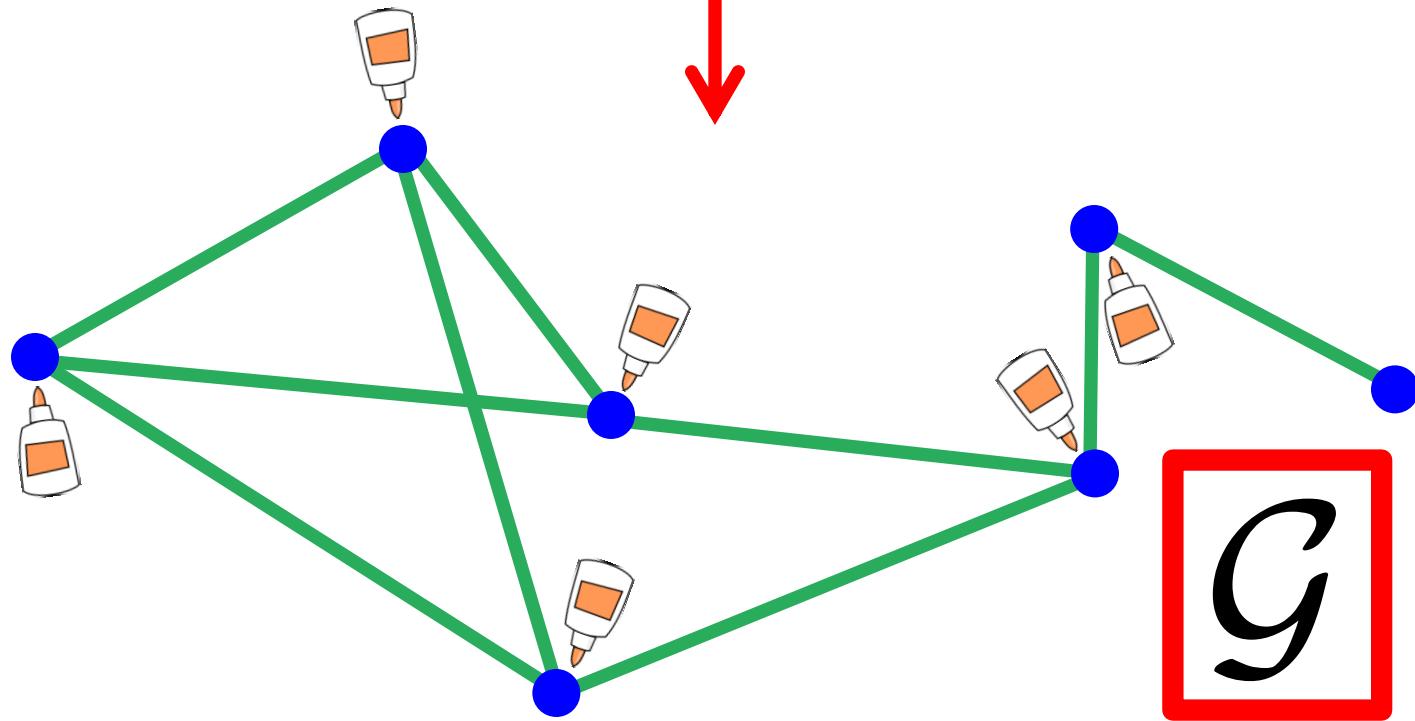
Soon-Yeong CHUNG*, Yun-Sung CHUNG** and Jong-Ho KIM***

The main goal of this paper is to study the discrete analog of the heat equation on graphs. The type of wave equation studied here is the edge-based Laplacian, which is a natural extension of the continuous Laplacian. The main result is that the solution of the heat equation on a graph can be expressed in terms of the eigenvalues and eigenvectors of the edge-based Laplacian. This result is analogous to the well-known fact that the solution of the heat equation on a domain in Euclidean space can be expressed in terms of the eigenvalues and eigenvectors of the Laplacian on that domain.

Traditionally, the study of wave equations on graphs has been limited to the case where the graph is a tree or a complete graph. In this paper, we extend the results to more general graphs, such as complete graphs and complete bipartite graphs. We also study the case where the graph is not connected, and show that the solution of the heat equation on such a graph can be expressed in terms of the eigenvalues and eigenvectors of the edge-based Laplacian on each connected component of the graph.

Geometric Realization

$$G = (V, E)$$



Edge Interior

$$\left(\begin{array}{c} v \in T\mathcal{G} \\ \text{---} \\ (a, b) \subset \mathbb{R}^1 \\ l_e = |b - a| \\ \text{---} \\ e \in E \end{array} \right)$$

Allowable operations:

$$\boxed{\nabla f}$$

$$\boxed{\nabla_{\text{calc}} \cdot X}$$

Admits differential structure

Two Measures

Discrete **vertex** measure

$$\mathcal{V}$$

Lebesgue **edge** measure

$$\mathcal{E}$$

Integrating Factors

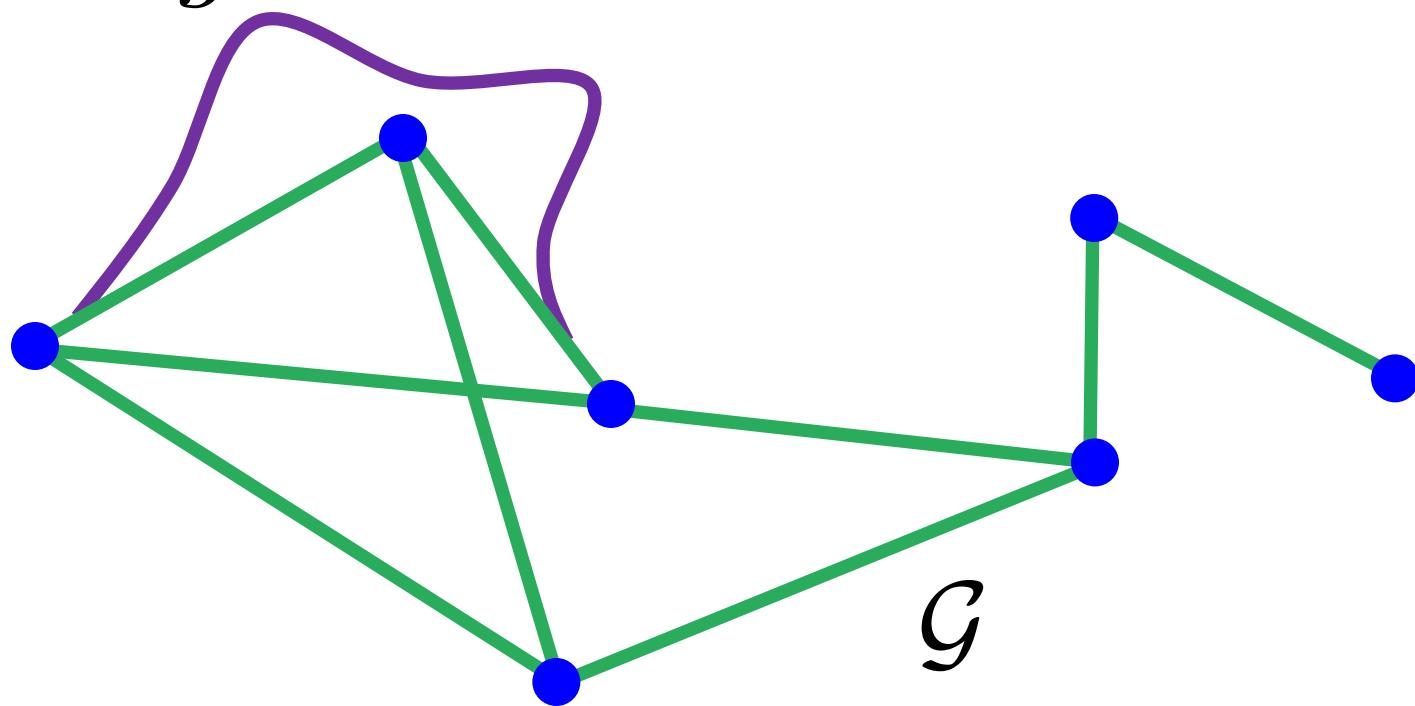
$$d\Gamma = \alpha \ d\mathcal{V} + \beta \ d\mathcal{E}$$

$$\int_{\mathcal{G}} f \ d\Gamma = \int_{\mathcal{G}} f\alpha \ d\mathcal{V} + \int_{\mathcal{G}} f\beta \ d\mathcal{E}$$

Integration by Parts

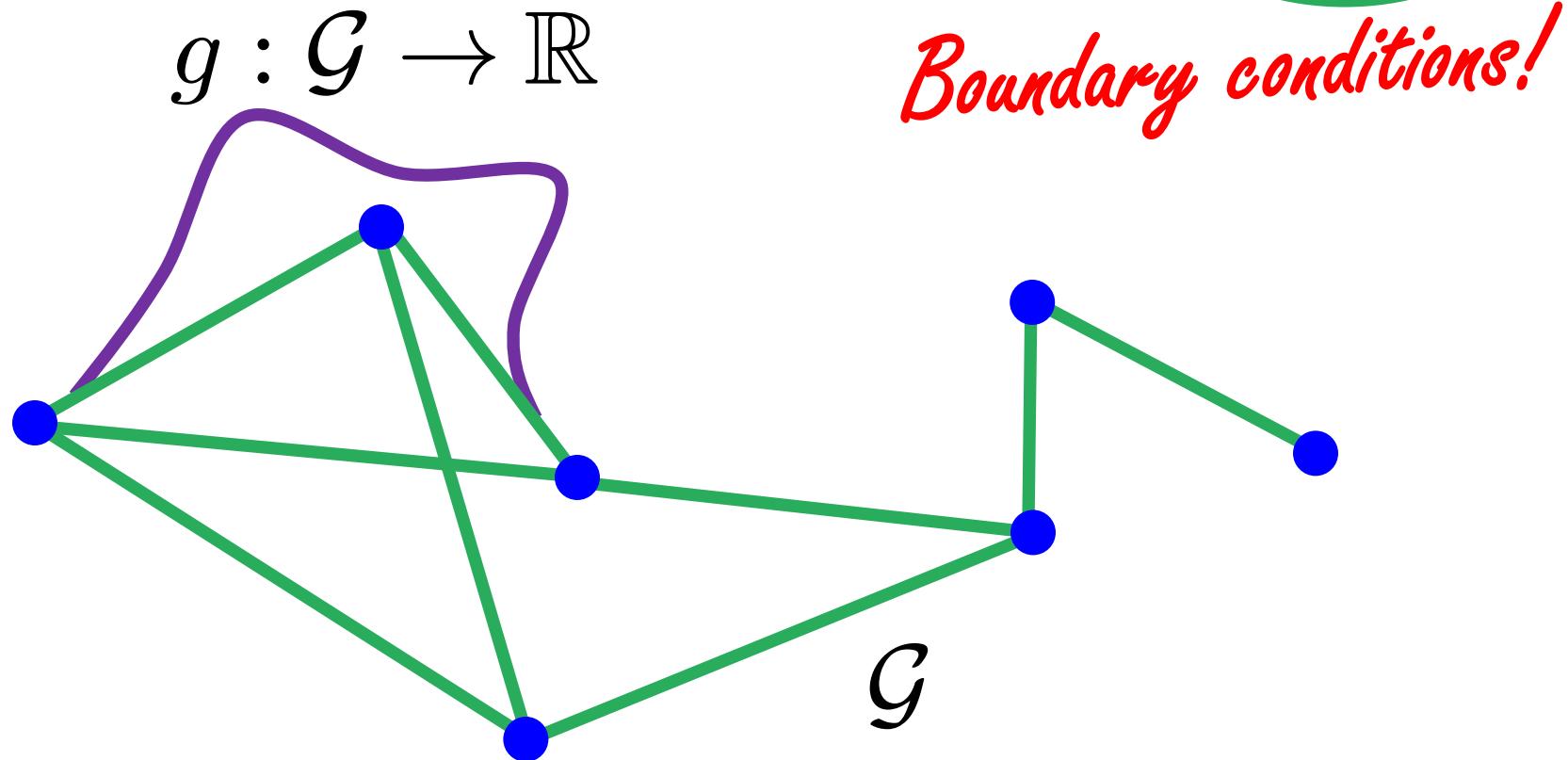
$$\mathcal{D}_X(g) \equiv - \int_{\mathcal{G} \setminus V} X \cdot \nabla g \ d\mathcal{E} = \int_{\mathcal{G}} [(\nabla_{calc} \cdot X)g \ d\mathcal{E} - (\tilde{n} \cdot X)g \ dV]$$

$$g : \mathcal{G} \rightarrow \mathbb{R}$$



Integration by Parts

$$\mathcal{D}_X(g) \equiv - \int_{\mathcal{G} \setminus V} X \cdot \nabla g \, d\mathcal{E} = \int_{\mathcal{G}} [(\nabla_{calc} \cdot X)g \, d\mathcal{E} - (\tilde{n} \cdot X)g \, dV]$$

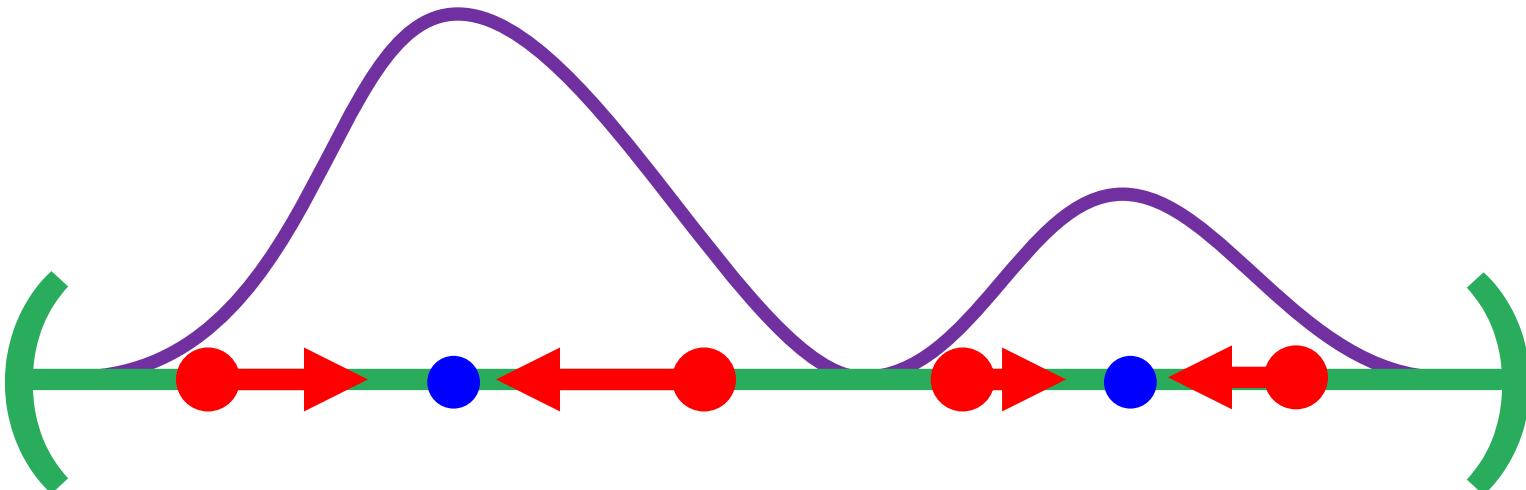


Divergence Factor

$$d\mathcal{D}_X \equiv (\nabla_{calc} \cdot X) d\mathcal{E} - (\tilde{n} \cdot X) d\mathcal{V}$$

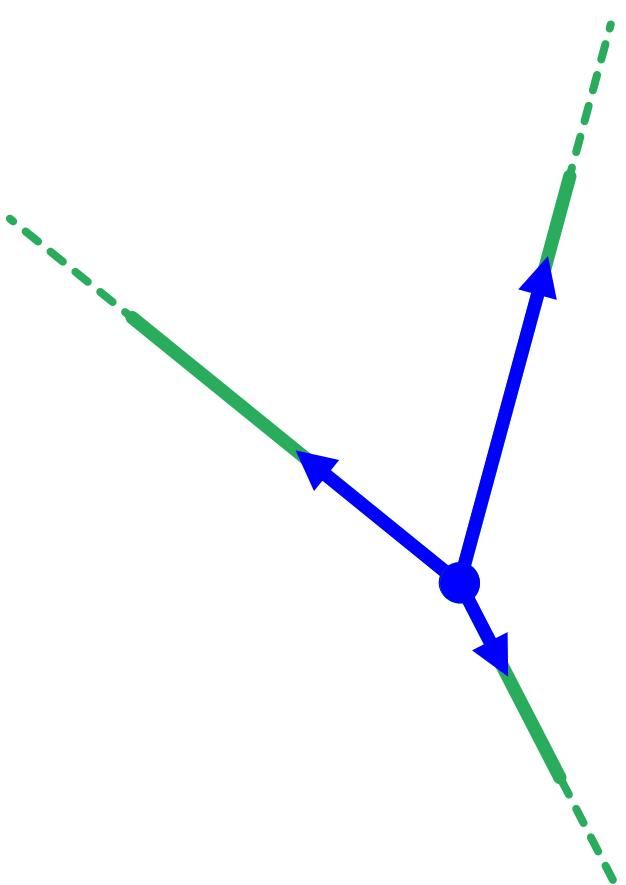


Divergence in edge interior



Divergence Factor

$$d\mathcal{D}_X \equiv (\nabla_{calc} \cdot X) d\mathcal{E} - (\tilde{n} \cdot X) d\mathcal{V}$$



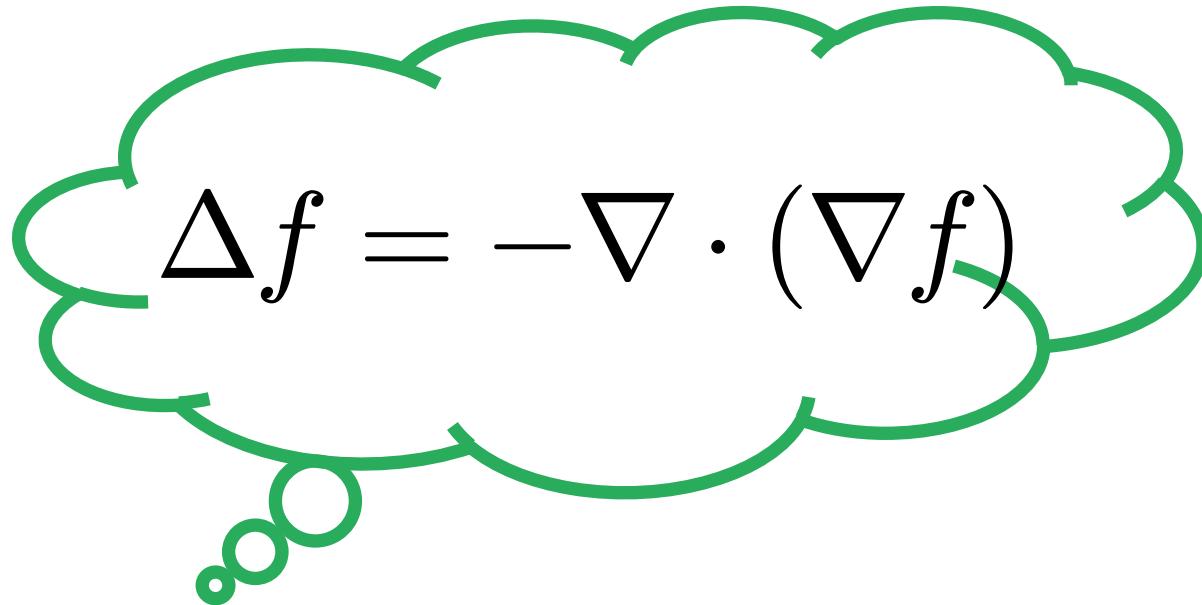
Divergence at vertices

Divergence Factor

$$d\mathcal{L}_f \equiv -d\mathcal{D}_{\nabla f}$$

Divergence Factor

$$d\mathcal{L}_f \equiv -d\mathcal{D}_{\nabla f}$$



Divergence Factor

$$d\mathcal{L}_f \equiv -d\mathcal{D}_{\nabla f}$$

$$d\mathcal{L}_f = \Delta_E d\mathcal{E} + \Delta_V d\mathcal{V}$$

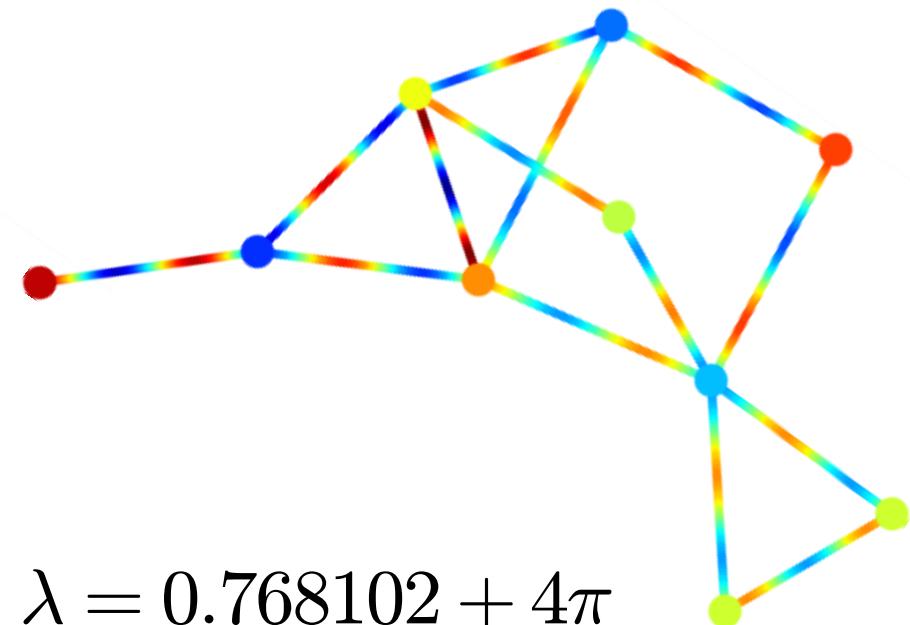
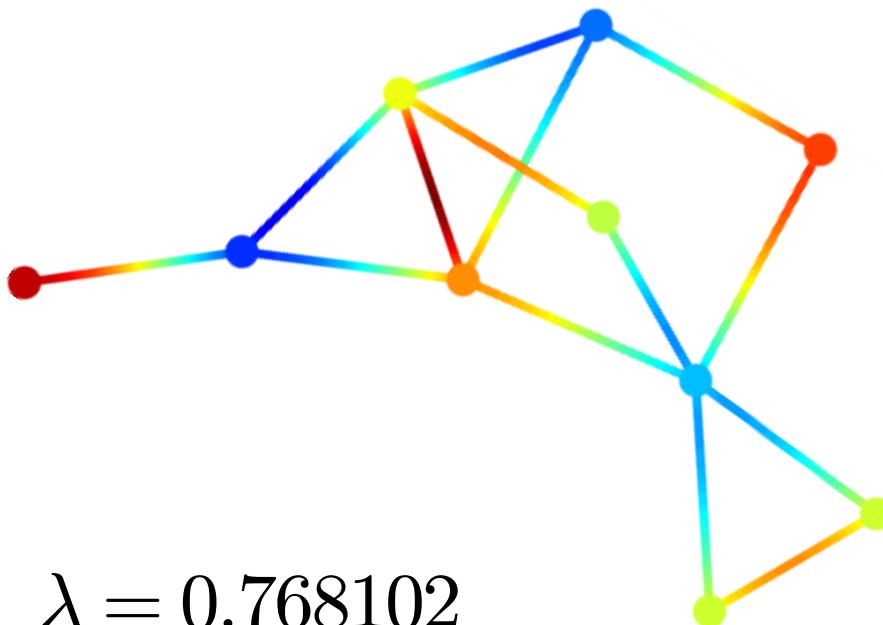
$$\Delta_E f = -\nabla_{calc} \cdot \nabla f$$

$$\Delta_V f = \tilde{n} \cdot \nabla f$$

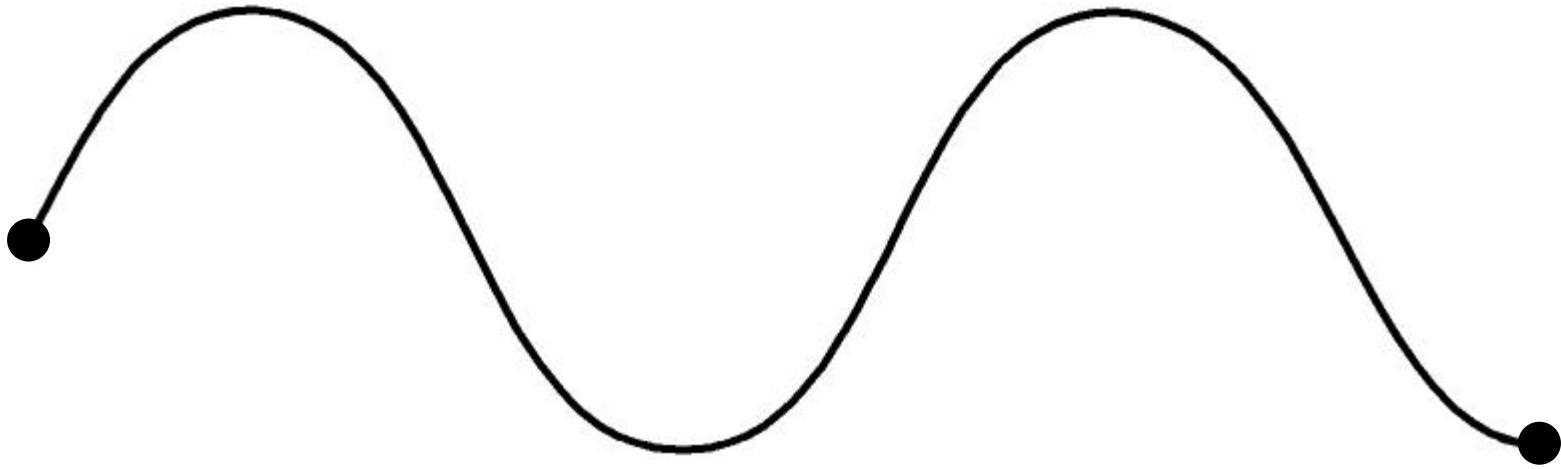
Edge-Based Eigenfunctions

$$\Delta_E f = -\nabla_{calc} \cdot \nabla f = \lambda f$$

$$\Delta_V f = \tilde{n} \cdot \nabla f = 0$$



Eigenfunction in Edge Interior



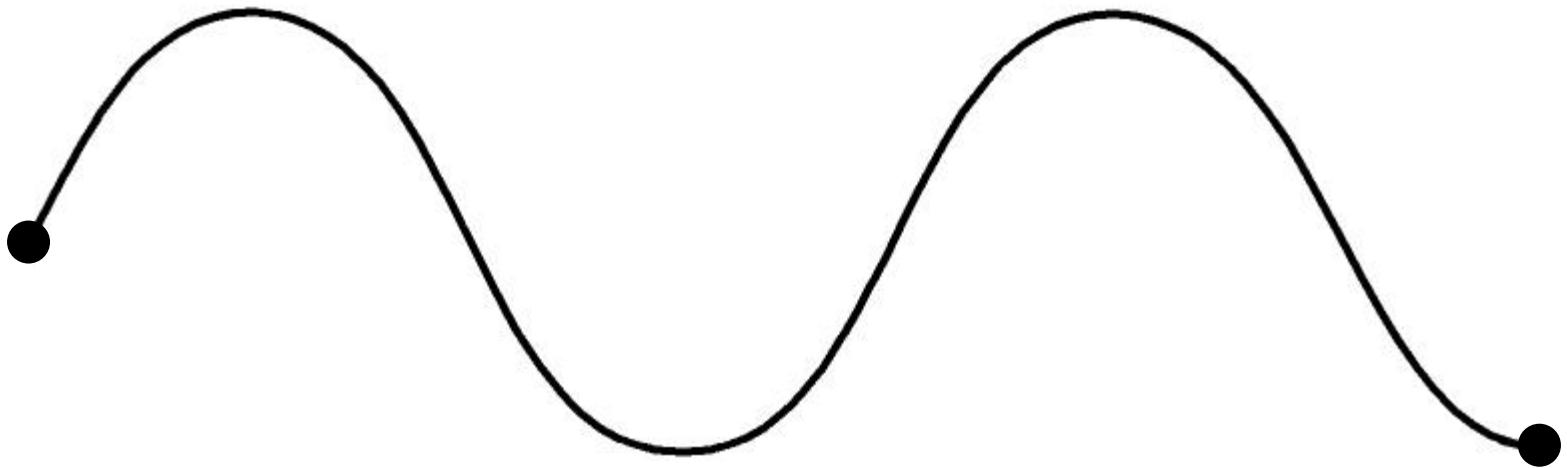
$$f(u) = f_e(0) = A \cos B$$

$$f(v) = f_e(l_e) = A \cos(\omega l_e + B) = A \cos(\omega l_e) \cos(B) - A \sin(\omega l_e) \sin(B)$$

$$f'(0) = -A\omega \sin B = -\omega \frac{f(v) - \cos(\omega l_e)f(u)}{\sin(\omega l_e)}$$

$$0 = \Delta_V f|_u = \sum_{e=(u,v) \in E} \frac{f(v) - \cos(\omega l_e)f(u)}{\sin(\omega l_e)}$$

Eigenfunction in Edge Interior



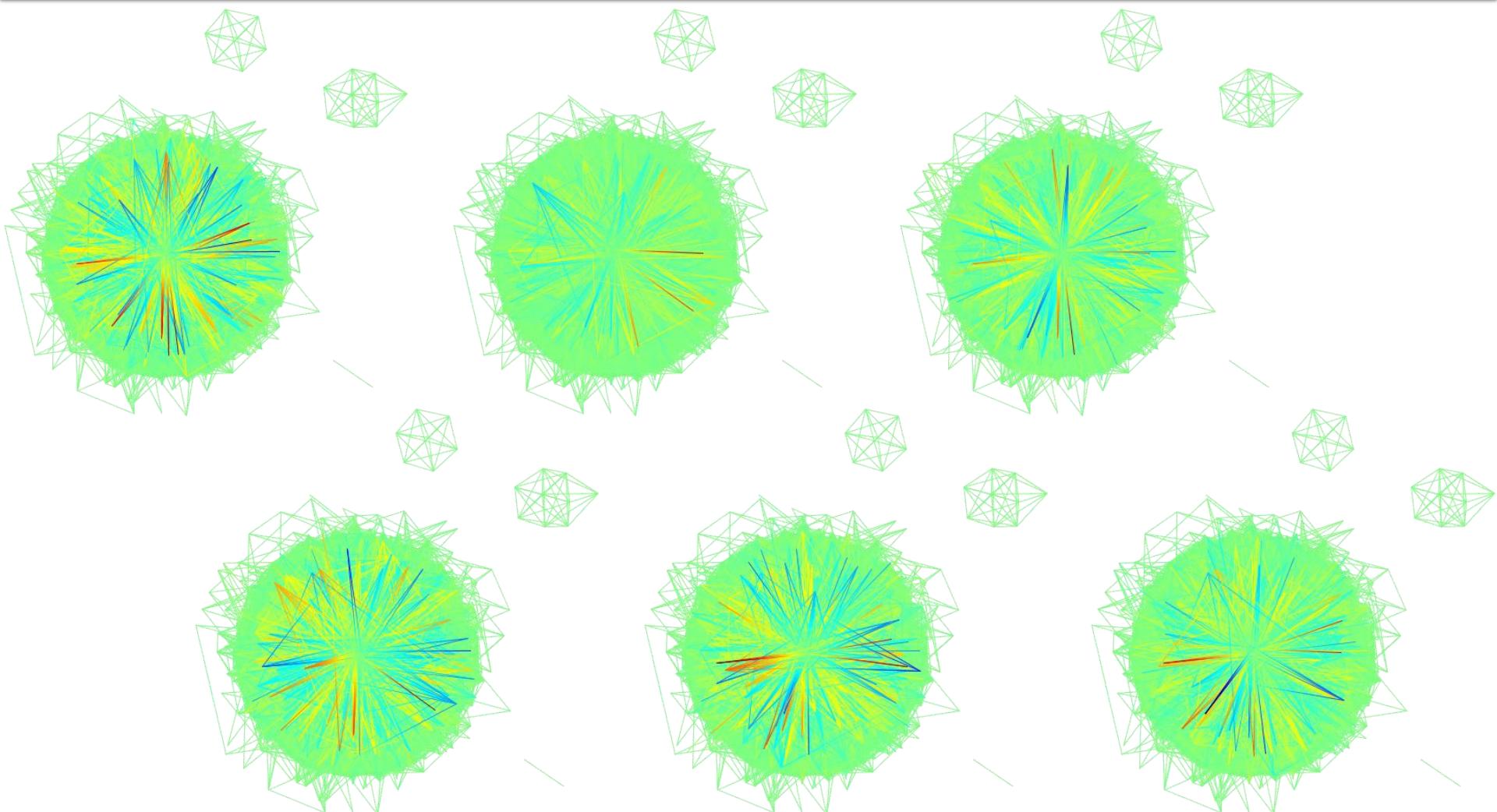
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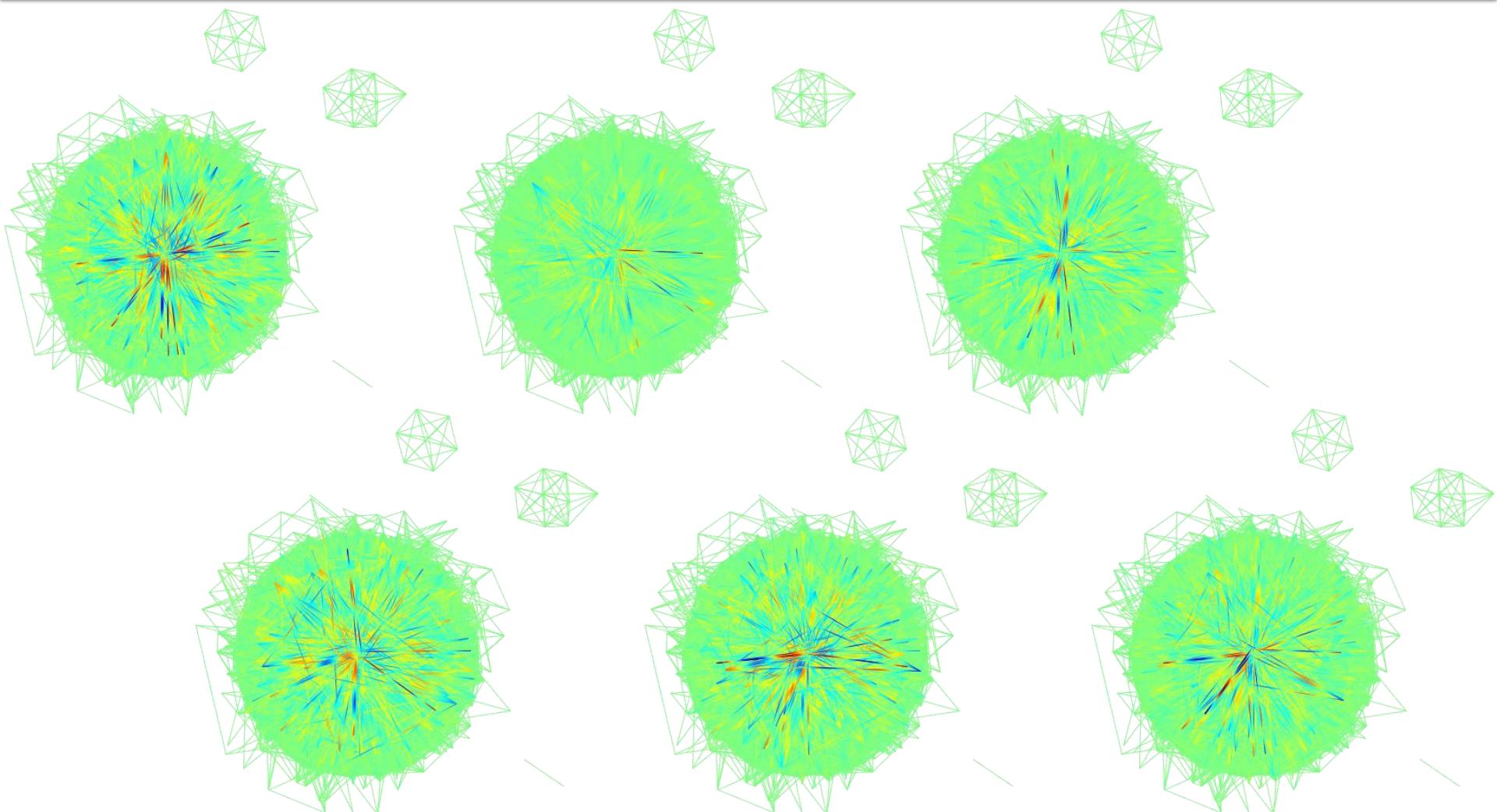
$$0 = \Delta_V f|_u = \sum_{e=(u,v) \in E} \frac{f(v) - \cos(\omega l_e)f(u)}{\sin(\omega l_e)}$$

Edge-Based Eigenfunctions



"Marvel Universe Looks Almost Like a Social Network" / Alberich, Miro-Julia, Rossello

Edge-Based Eigenfunctions



"Marvel Universe Looks Almost Like a Social Network" / Alberich, Miro-Julia, Rossello

Geometric Graph Wave Equation

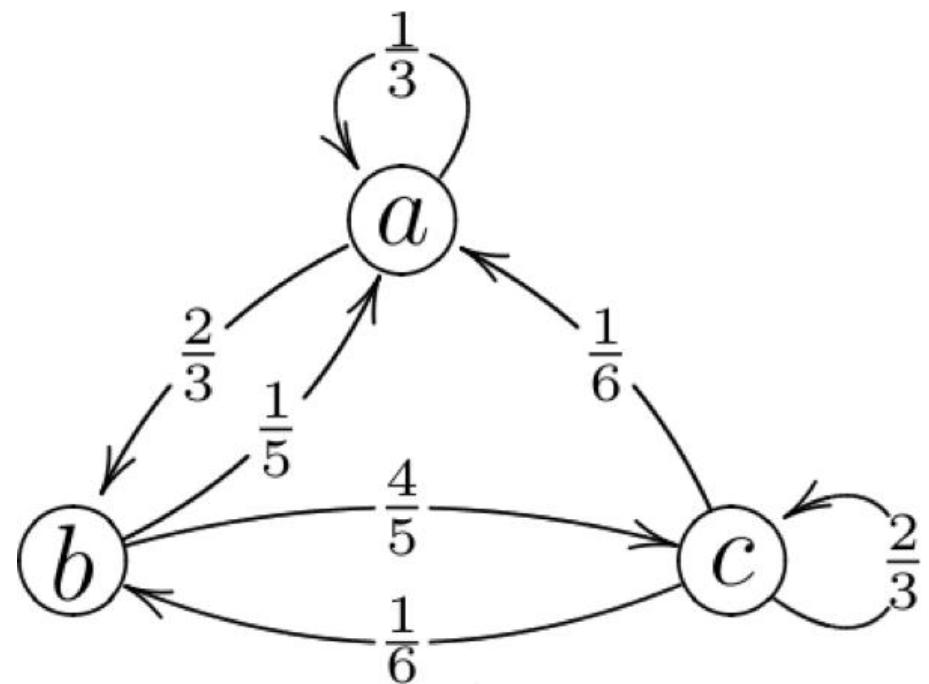
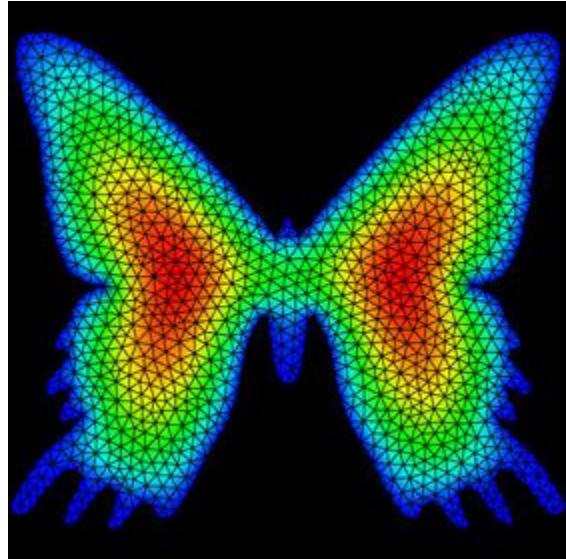
$$(\alpha d\mathcal{V} + \beta d\mathcal{E})u_{tt} = d\mathcal{D}_\gamma \nabla u$$

- Encodes boundary conditions
- Preserves energy
- Constant speed of propagation

Geometrization: Conclusion

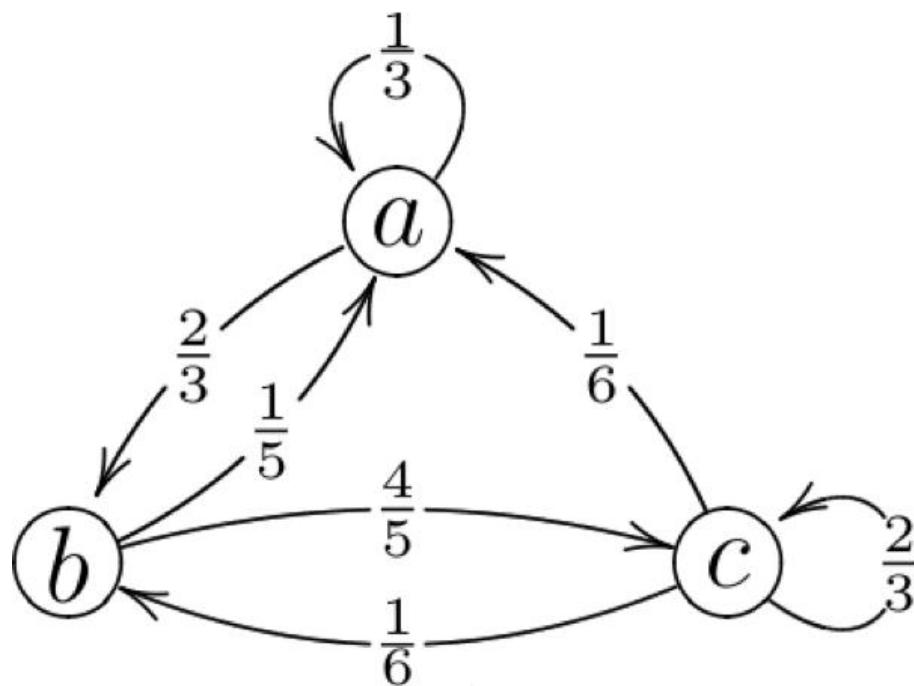
Mostly theoretical, needs
more work for practical
applications.

Discrete Laplacians



Markov Chain Stationary States

$$\Delta(u, v) \equiv [I - P](u, v) = \begin{cases} 1 & \text{if } u = v \\ -P(u, v) & \text{if } u \text{ and } v \text{ are adjacent} \\ 0 & \text{otherwise} \end{cases}$$



Markov Chain Stationary States

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$$\Delta(c_1 u + c_2 v) = c_1 \Delta u + c_2 \Delta v$$

- **Compact and bounded**

$$\|\Delta u\| \leq M \|u\| \quad \forall u$$

- **Self-adjoint**

$$\langle \Delta u, v \rangle = \langle u, \Delta v \rangle$$

- **Categorizes extrema**

$$x \in \Omega \text{ local minimum} \implies [\Delta u](x) \leq 0$$

Markov Chain Stationary States

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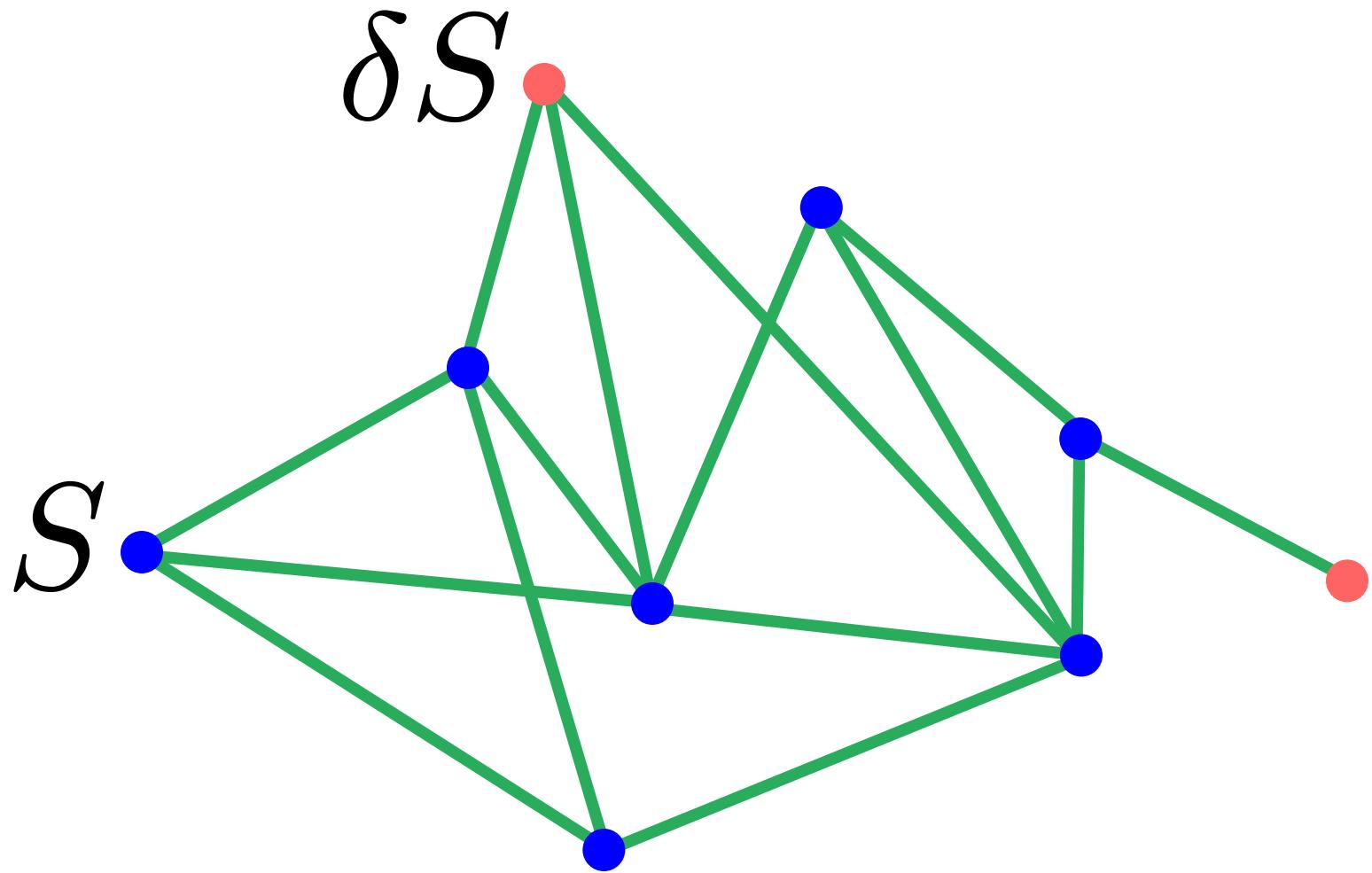
$$\langle \Delta u, v \rangle = \langle u, \Delta v \rangle$$

$$\mathcal{L} = T^{1/2} \Delta T^{-1/2}$$

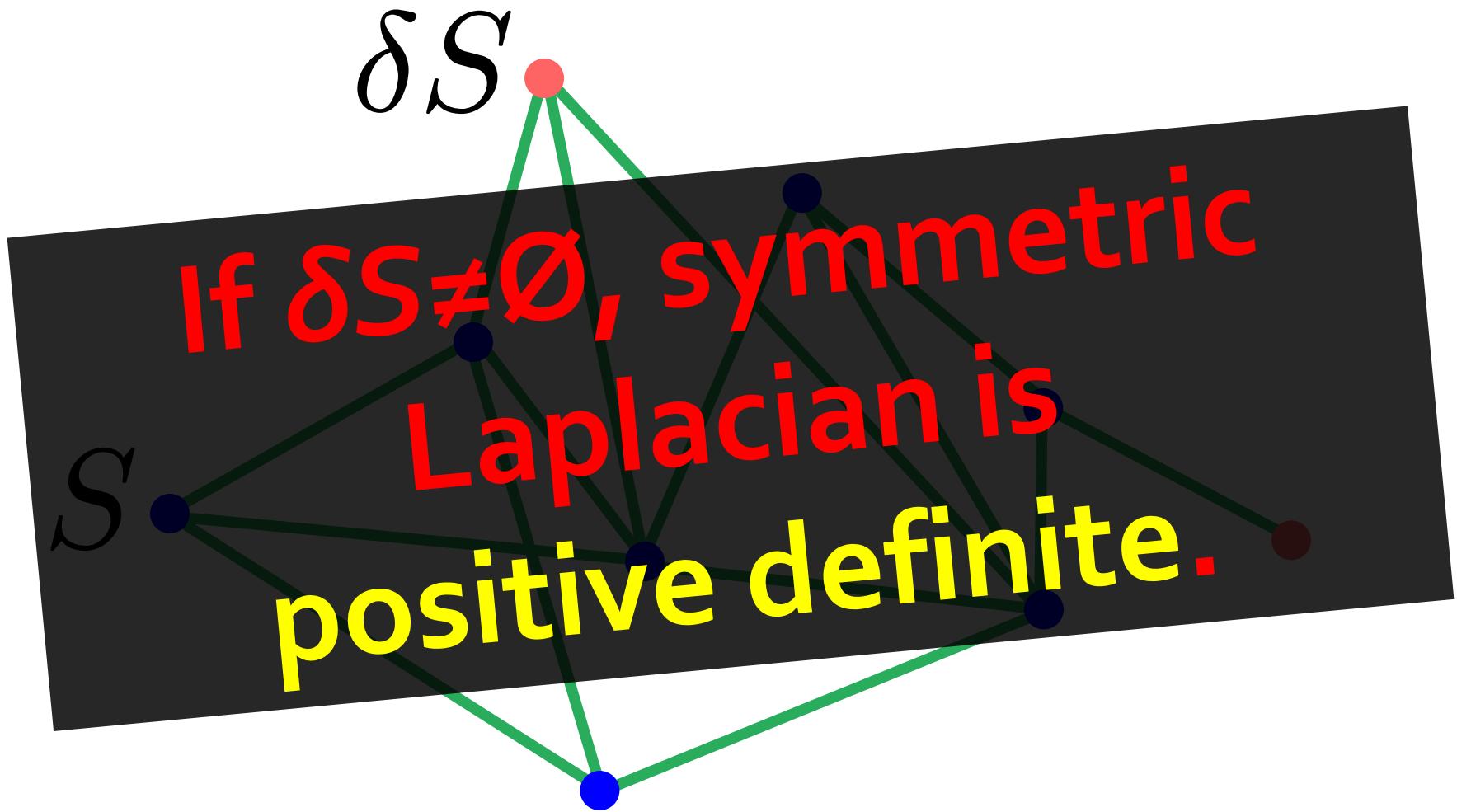
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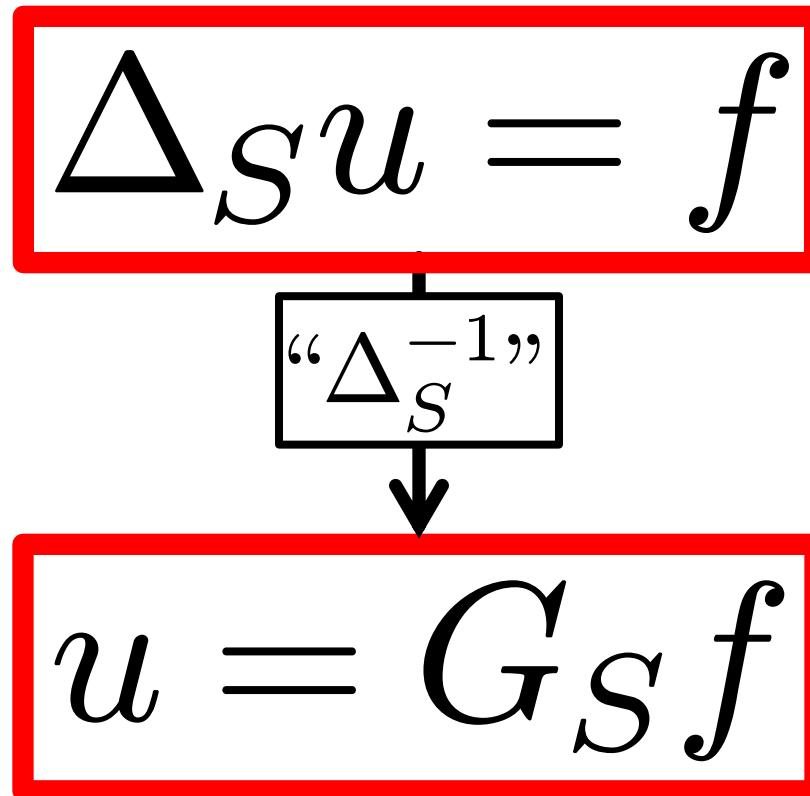
Discrete Setup



Discrete Setup



Discrete Green's Functions



Closed-Form Laplace Solution

Lemma 7 ([CY00], Theorem 1). Suppose $\Delta f = 0$ on S with $f(x) = \sigma(x)$ for $x \in \delta S$. Then, f satisfies

$$f(z) = d_z^{-1/2} \sum_i \frac{1}{\lambda_i} \phi_i(z) \sum_{\substack{x \in S \\ (x,y) \in E \\ y \in \delta S}} d_x^{-1/2} \phi_i(x) \sigma(y)$$

$$\mathcal{L} = T^{1/2} \Delta T^{-1/2}$$

$$\mathcal{L} \phi_i = \lambda_i \phi_i$$

Proof strategy: Construct solution to symmetric Laplace equation from f , expand its projection onto Fourier basis.

Semi-Discrete PDEs

- Discrete along the graph
- Continuous in time

$$u_t = -\Delta u$$

$$u_{tt} = -\Delta u$$

Two Dual Viewpoints

$$u_t = -\Delta u$$

Heat equation

$$u_{tt} = -\Delta u$$

Wave equation

PDE standpoint

Two Dual Viewpoints

u_t Properties similar to
physical phenomena.

Δu

$u_{tt} = -\Delta u$

Heat equation Wave equation

PDE standpoint

Two Dual Viewpoints

$$\vec{u}_t = M \vec{u}$$

Linear ordinary
differential equation

ODE standpoint

Two Dual Viewpoints

Can solve without convergence arguments
at Linear ordinary differential equation
and other technicalities.

ODE standpoint

Huygens Property

Lemma 11 ([CCK07], “Huygens Property” Theorem 3.6). *Suppose u satisfies the heat equation $u_t = -\Delta u$ on a graph with $S = V$. Then, for every $t, \delta > 0$ we have*

$$u(x, t + \delta) = \langle K(x, \cdot, \delta), u(\cdot, t) \rangle \quad (54)$$

Huygens Property

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$$u(x, t + \delta) = \langle K(x, \cdot, \delta), u(\cdot, t) \rangle \quad (54)$$

Heat kernel

$$u(x, t) = \langle K_S(x, \cdot, t), f(\cdot) \rangle_S$$

$$K_S(u, v, t) = \sum_{i=1}^{|S|} e^{-\lambda_i t} \phi_i(u) \phi_i(v) \sqrt{\frac{d_v}{d_u}}$$

Huygens Property

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$$u(x, t + \delta) = \langle K(x, \cdot, \delta), u(\cdot, t) \rangle \quad (54)$$

$$u(x, t) = \frac{1}{\sqrt{d_x}} \sum_i c_i e^{-\lambda_i t} \phi_i(x)$$

Huygens Property

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$$\begin{aligned} u(x, t + \delta) &= \frac{1}{\sqrt{d_x}} \sum_i c_i e^{-\lambda_i(t+\delta)} \phi_i(x) & u(x, t) &= \frac{1}{\sqrt{d_x}} \sum_i c_i e^{-\lambda_i t} \phi_i(x) \\ &= \frac{1}{\sqrt{d_x}} \sum_i e^{-\lambda_i \delta} c_i e^{-\lambda_i t} \phi_i(x) \end{aligned}$$

Huygens Property

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$$\begin{aligned} u(x, t + \delta) &= \frac{1}{\sqrt{d_x}} \sum_i c_i e^{-\lambda_i(t+\delta)} \phi_i(x) & u(x, t) &= \frac{1}{\sqrt{d_x}} \sum_i c_i e^{-\lambda_i t} \phi_i(x) \\ &= \frac{1}{\sqrt{d_x}} \sum_i e^{-\lambda_i \delta} c_i e^{-\lambda_i t} \phi_i(x) \\ &= \frac{1}{\sqrt{d_x}} \sum_i e^{-\lambda_i \delta} \phi_i(x) \left\langle \sqrt{d_y} \phi_i(y), \frac{1}{\sqrt{d_y}} \sum_j c_j e^{-\lambda_j t} \phi_j(y) \right\rangle \end{aligned}$$

Huygens Property

Lemma 11 ([CCK07], “Huygens Property” Theorem 3.6). Suppose u satisfies the heat equation $u_t = -\Delta u$ on a graph with $S = V$. Then, for every $t, \delta > 0$ we have

$$u(x, t + \delta) = \langle K(x, \cdot, \delta), u(\cdot, t) \rangle \quad (54)$$

$$\begin{aligned} u(x, t + \delta) &= \frac{1}{\sqrt{d_x}} \sum_i c_i e^{-\lambda_i(t+\delta)} \phi_i(x) & u(x, t) &= \frac{1}{\sqrt{d_x}} \sum_i c_i e^{-\lambda_i t} \phi_i(x) \\ &= \frac{1}{\sqrt{d_x}} \sum_i e^{-\lambda_i \delta} c_i e^{-\lambda_i t} \phi_i(x) \\ &= \frac{1}{\sqrt{d_x}} \sum_i e^{-\lambda_i \delta} \phi_i(x) \left\langle \sqrt{d_y} \phi_i(y), \frac{1}{\sqrt{d_y}} \sum_j c_j e^{-\lambda_j t} \phi_j(y) \right\rangle \\ &= \frac{1}{\sqrt{d_x}} \sum_i e^{-\lambda_i \delta} \phi_i(x) \langle \sqrt{d_y} \phi_i(y), u(\cdot, t) \rangle \end{aligned}$$

Huygens Property

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$$u(x, t + \delta) = \frac{1}{\sqrt{d_x}} \sum_i c_i e^{-\lambda_i(t+\delta)} \phi_i(x) \quad u(x, t) = \frac{1}{\sqrt{d_x}} \sum_i c_i e^{-\lambda_i t} \phi_i(x)$$

Generic proof strategy

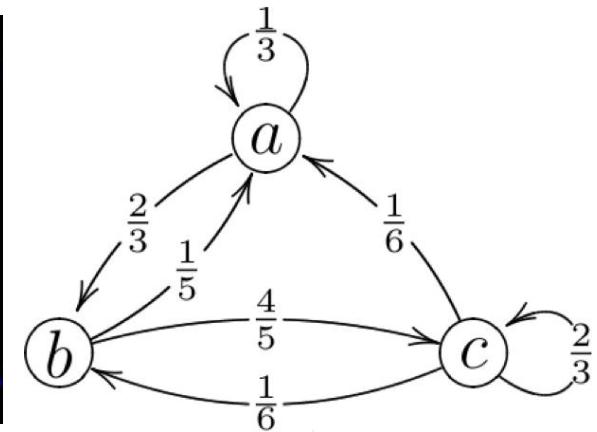
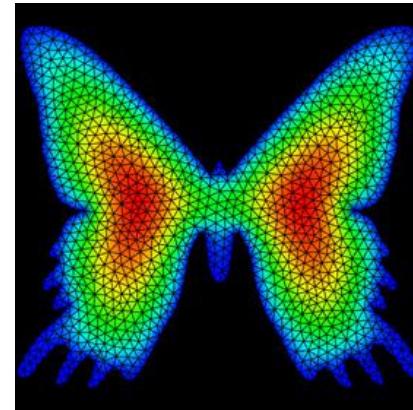
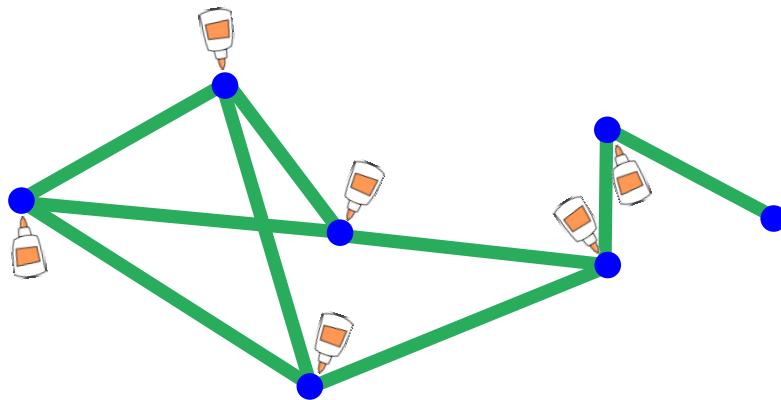
doesn't change from
continuous case.

$$\begin{aligned} &= \frac{1}{\sqrt{d_x}} \sum_i c_i e^{-\lambda_i \delta} \phi_i(x) \langle \sqrt{d_y} \phi_i(y), \frac{1}{\sqrt{d_y}} \sum_j c_j e^{-\lambda_j t} \phi_j(y) \rangle \\ &= \langle K(x, \cdot, \delta), u(\cdot, t) \rangle \end{aligned}$$

Themes

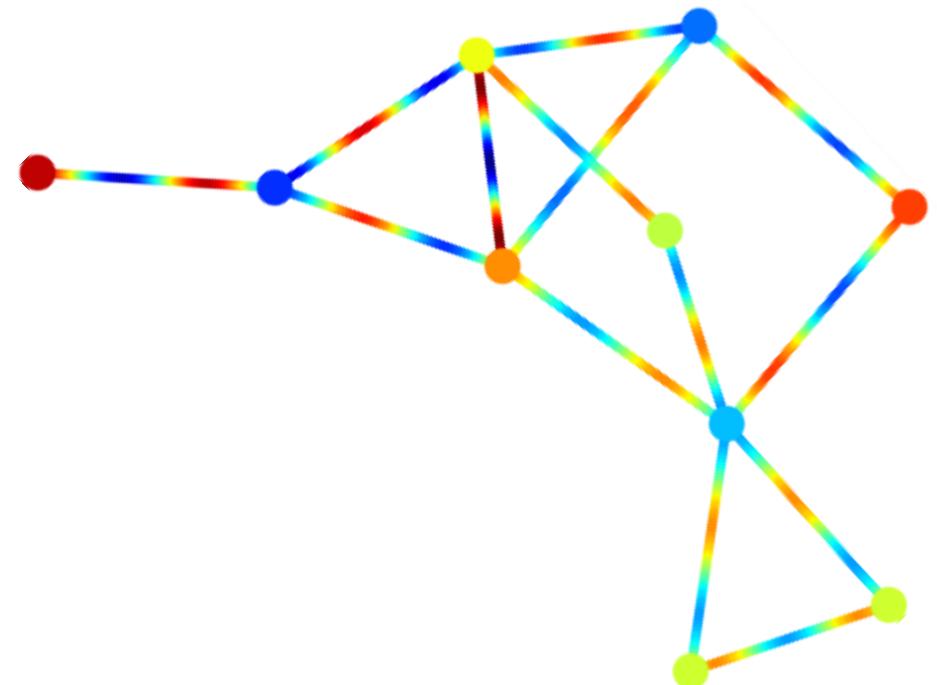
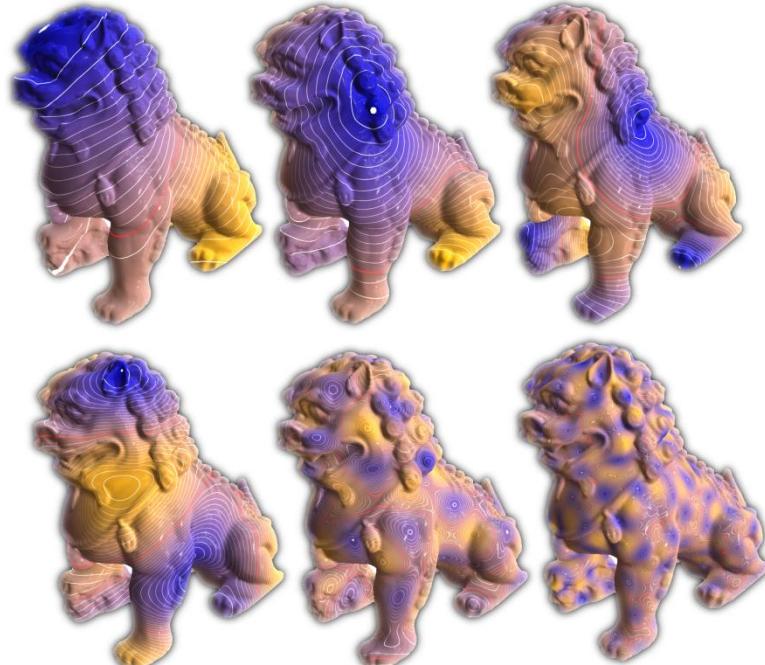
Ingredients:

1. Domain
2. “Differential” operator



Themes

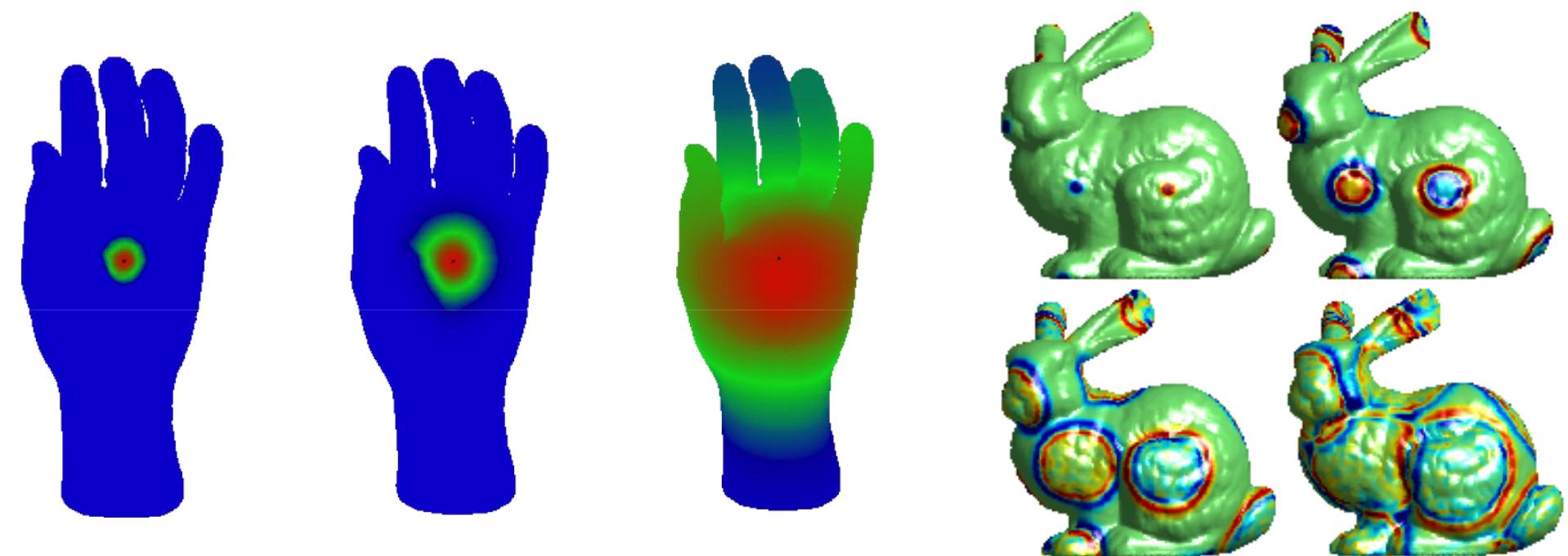
Hammer:
Eigen-analysis



Themes

Result:

PDE-like behavior

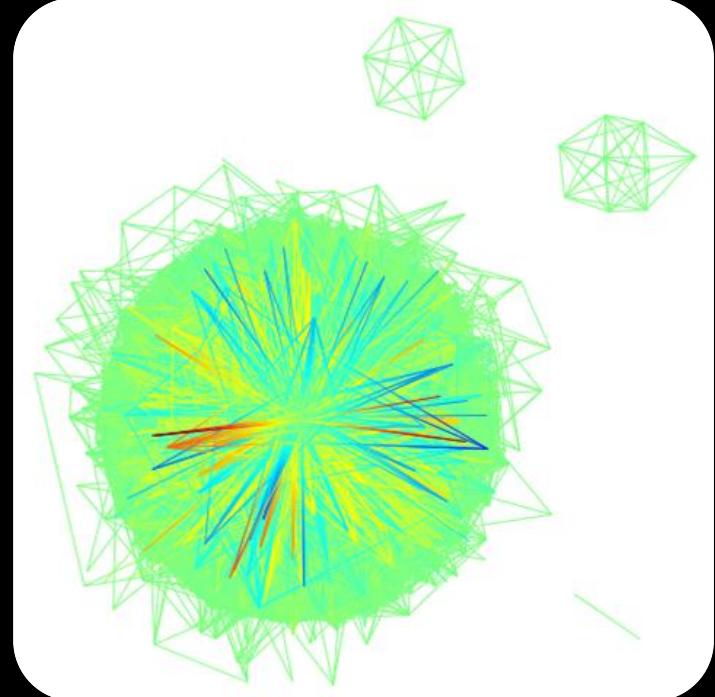
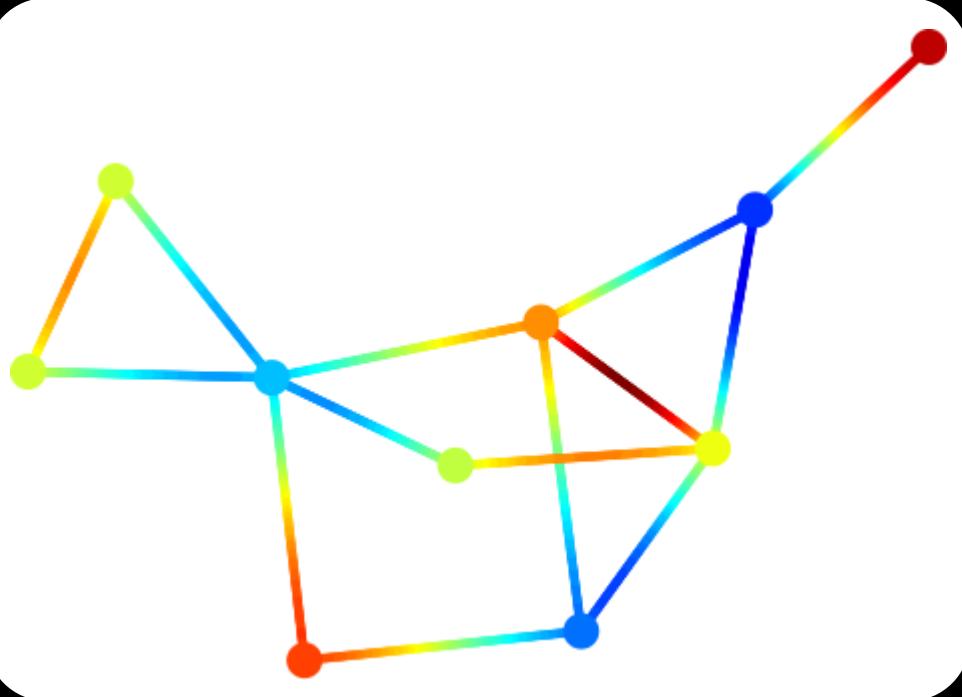


In Other Words...

PDE behavior can be
domain-independent
in a rigorous way.

Next Steps

- Improve **practicality** of edge-based approach
- Adaptation of **geometric methods** to graph problems
- Organized **advantages/limitations** of each representation



PDE Approaches to Graph Analysis

Questions?